



Workshop on CSR in Storage Rings

Napa, October 27-28, 2002

CSR and Longitudinal Beam Dynamics: a Numerical Study

M. Venturini,

SLAC

(with R. Warnock and R. Ruth)

Outline

- 1D Vlasov-Fokker-Planck equation to study interplay between CSR and its dynamical effects.
- Simplified model for CSR impedance (parallel plates)
- Solve VFP equation numerically.
- Two regimes:
 - » Bursting (**NSLS VUV Storage Ring**)
 - » Steady state emissions (**Bessy II**).

In Vlasov (Vino?) Veritas

VFP Equation

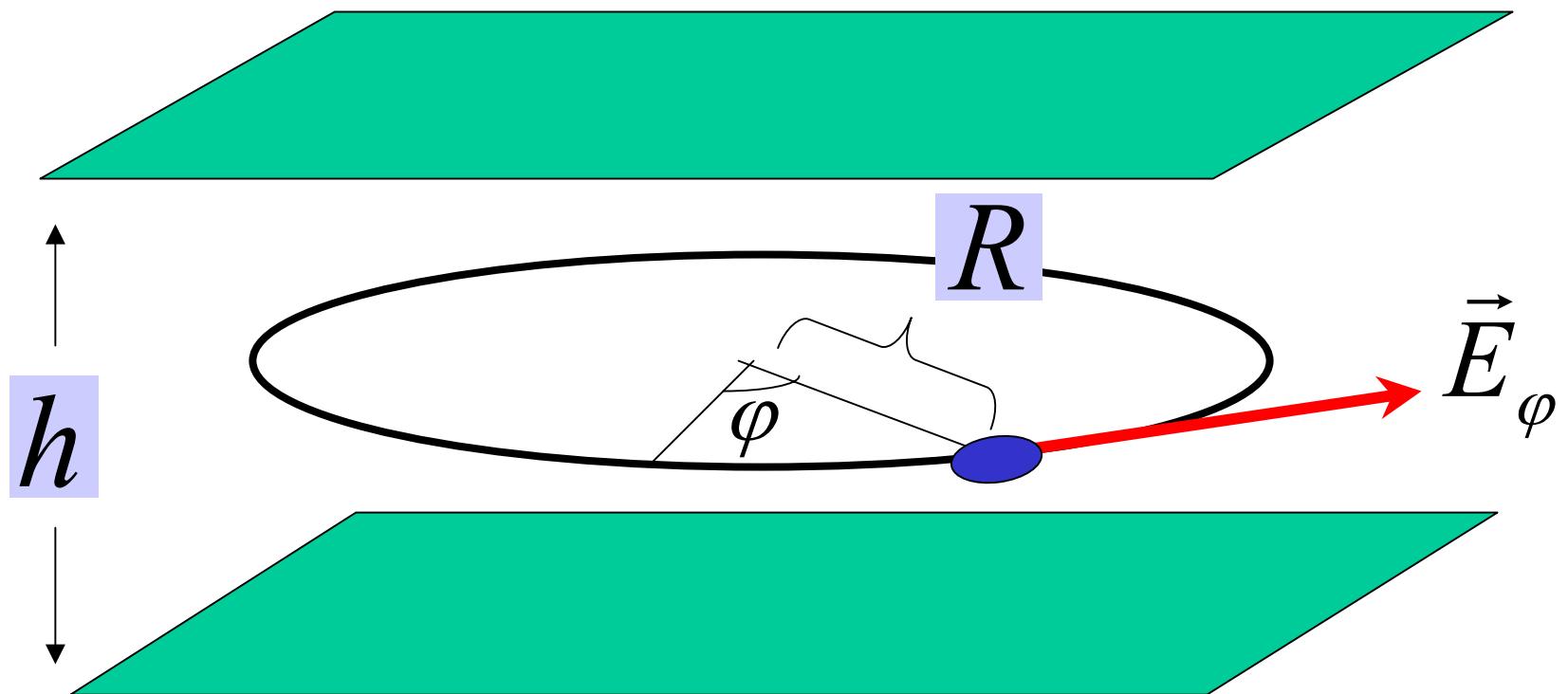
$$q = z / \sigma_z, \quad p = -\Delta E / \sigma_E$$

$$\frac{\partial f}{\partial \tau} + p \frac{\partial f}{\partial q} - [q + F_c(q, \tau, f)] \frac{\partial f}{\partial p} = \frac{2}{\omega_s t_d} \frac{\partial}{\partial p} \left(p f + \frac{\partial f}{\partial p} \right)$$

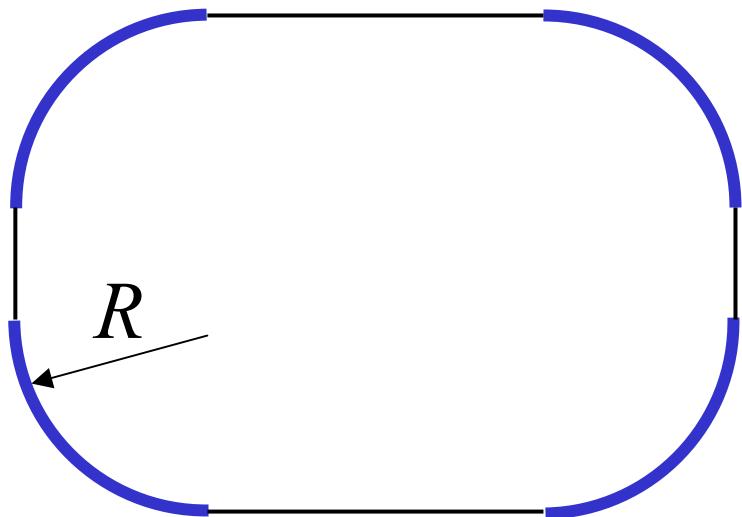
Annotations pointing to specific terms in the equation:

- $\tau = \omega_s t$ points to $\frac{\partial f}{\partial \tau}$
- "RF focusing" points to $p \frac{\partial f}{\partial q}$
- "Collective Force" points to $[q + F_c(q, \tau, f)] \frac{\partial f}{\partial p}$
- "Damping" points to $\frac{2}{\omega_s t_d} \frac{\partial}{\partial p} \left(p f + \frac{\partial f}{\partial p} \right)$
- "Quantum Excitations" points to $\frac{\partial f}{\partial \tau}$

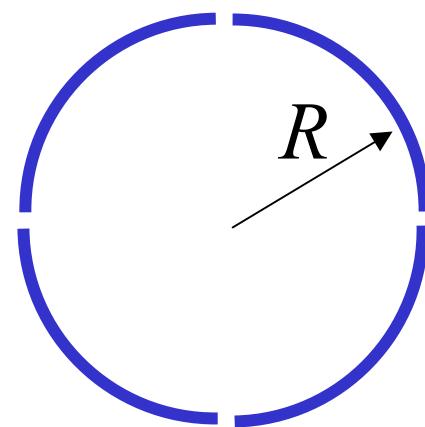
Parallel Plate Model for CSR: Geometry Outline



Actual lattice



Simplified lattice for
CSR force calculation



$$2\pi \langle R \rangle E_\varphi^{\langle R \rangle} = 2\pi R E_\varphi^R$$

Collective Force due CSR

(also, see R. Warnock talk)

FT of (normalized) charge density of bunch.

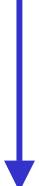
$$E_\varphi = -\frac{e\omega_0 N}{(2\pi)^2 R} \sum_n e^{in\varphi} \int_{-\infty}^{\infty} d\omega e^{-i\omega t} Z(n, \omega) \int_{-\infty}^t dt' e^{i(\omega - n\omega_0)t'} \lambda_n(t')$$

Assume charge distribution doesn't change much over one turn (rigid bunch approx).

$$E_\varphi = -\frac{e\omega_0 N}{2\pi R} \sum_n e^{in(\varphi - n\omega_0 t)} Z(n, n\omega_0) \lambda_n(t')$$

...more on collective force

$$\frac{\partial f}{\partial \tau} + p \frac{\partial f}{\partial q} - [q + F_c(q, \tau, f)] \frac{\partial f}{\partial p} = \frac{2}{\omega_s t_d} \frac{\partial}{\partial p} \left(p f + \frac{\partial f}{\partial p} \right)$$



$$F_c(q, \tau, f) = -I_c \omega_0 \sum_n e^{inq\sigma_z/R} Z(n) \lambda_n$$



$$I_c = \frac{e^2 N}{2\pi\nu_s \sigma_E} = \frac{e^2 N \sigma_z}{2\pi\alpha \langle R \rangle E_0 \sigma_\delta^2} \quad \left[\frac{C}{V} \right]$$

NSLS VUV Ring Parameters



Energy 737 MeV

Average machine radius 8.1 m

Local radius of curvature 1.9 m

Vacuum chamber **aperture** 4.2 cm

Nominal bunch length (rms) 5 cm

Nominal energy spread (rms) $\sigma_\delta = 5 \times 10^{-4}$

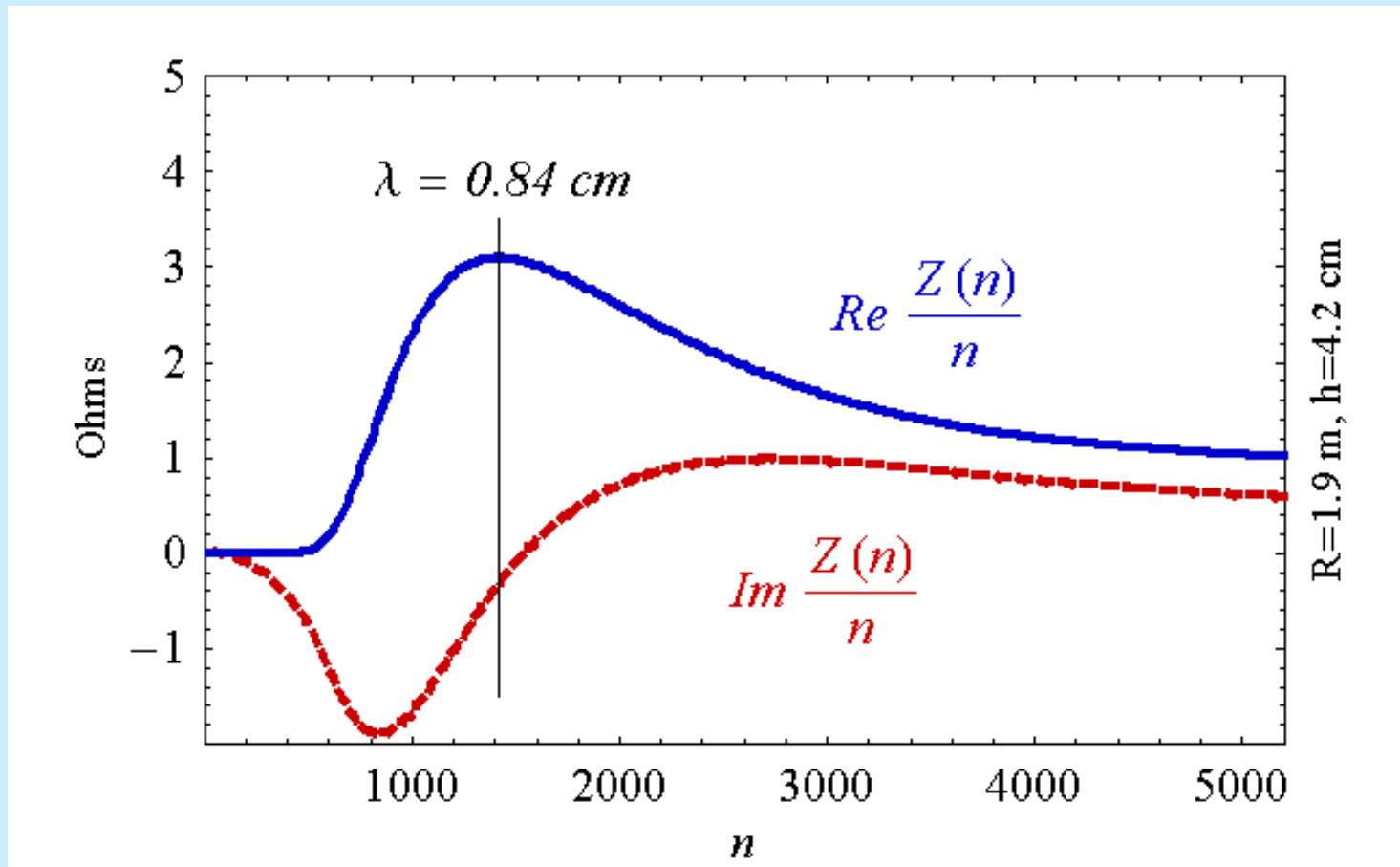
Synchrotron tune $\nu_s = 0.018$

Longitudinal damping time 10 ms



Parallel Plate Model : NSLS VUV Ring

$$Z(n) \equiv Z(n, n\omega_0)$$



Analytic Expression for CSR Impedance (Parallel Plates)

By definition: $-2\pi R E_\varphi(n, \omega) = Z(n, \omega) I_\varphi(n, \omega)$

Impedance

$$\frac{Z(n, \omega)}{n} = \frac{2\pi^2 Z_0}{\beta} \frac{R}{h} \sum_{p=1,3,\dots} \Lambda_p \left[\frac{\beta\omega R}{nc} J'_n H_n'^{(1)} + \frac{\alpha_p^2}{\gamma_p^2} J_n H_n^{(1)} \right]$$

With $\alpha_p = \pi p / h$, $\gamma_p^2 = (\omega/c)^2 - \alpha_p^2$

Argument of Bessel functions $R\gamma_p$

$\Lambda_p = \sin(x) / x$, $x = \pi p \delta h / 2h$

δh = beam height

Divide and Conquer

Numerical Method

- Solve VFP by operator splitting
(propagate distribution under Vlasov and Fokker-Planck parts separately).
- Propagate distribution function under FP by a simple Euler step using finite differences
(no need to be fancy here).

Propagation of Distribution under Vlasov

- Idea: follow the **particle trajectories** (method known as Semi-Lagrangian, Perron-Frobenius, Time-Splitting ...)
- Implemented with the Yabe *et al.* variation:
propagate derivatives $\partial_q f$, $\partial_{\bar{q}} f$ alongside with f

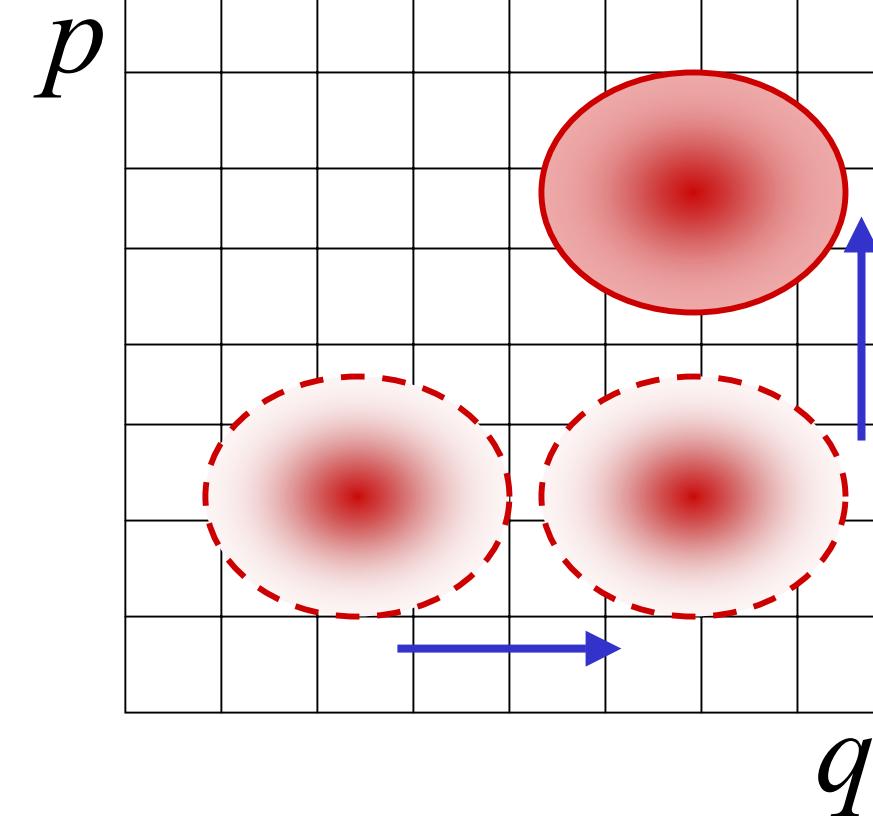
Perron-Frobenius Operator

- $z = (q, p)$ is point in phase-space.
- Mapping induced by the Hamiltonian dynamics from time t_0 to time t :

$$z \rightarrow z' = M_{t_0 \rightarrow t}(z)$$

- Propagation of f (Perron-Frobenius operator)

$$f(z, t) = f(M_{t_0 \rightarrow t}^{-1}(z), t_0)$$



- Do one time step along q , and one time step along p .

- Discretize in time; represent f on grid

$$f(z_{ij}, t + \Delta t) = f(M_{t \rightarrow t+\Delta t}^{-1}(z_{ij}), t_0)$$

- Break map into sequence of kicks

$$M_{t \rightarrow t+\Delta t} \approx M_{t \rightarrow t+\Delta t}^{(p)} \circ M_{t \rightarrow t+\Delta t}^{(q)}$$

\downarrow

$$M_{t \rightarrow t+\Delta t}^{(q)} : \begin{cases} p' = p \\ q' = q + p\Delta t \end{cases}$$

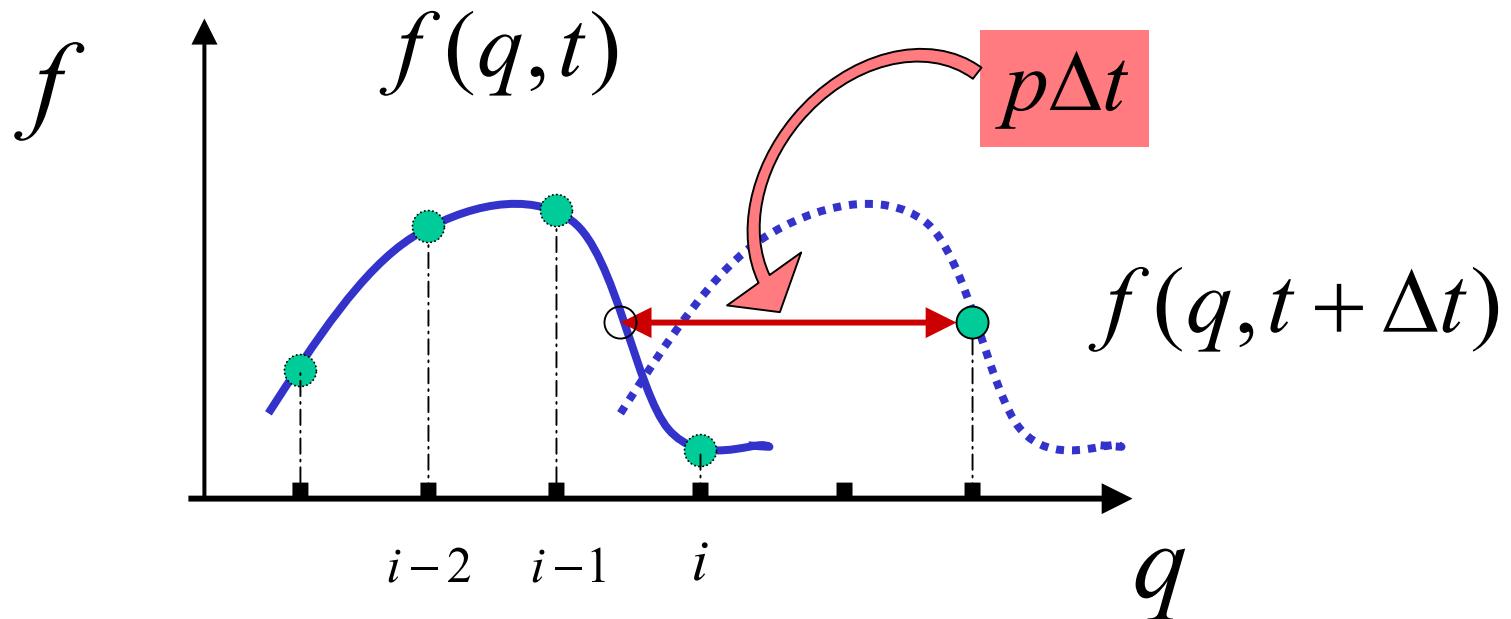
\downarrow

$$M_{t \rightarrow t+\Delta t}^{(p)} : \begin{cases} p' = p - [q + F_c(q, f)]\Delta t \\ q' = q \end{cases}$$

Step along q (pure drift)

$$f(q, t + \Delta t) = f(q - p\Delta t, t)$$

Do cubic interpolation between grid-points
Using f and $\partial_q f$ at end-points.



Propagation of Derivatives

$$\partial_q f(q, t + \Delta t) = \partial_q f(q - p\Delta t, t)$$

- ▶ Propagate $\partial_q f$ by differentiating same interpolating polynomial used to propagate f
-

- ▶ For $\partial_p f$ use finite differences. Start from Vlasov:

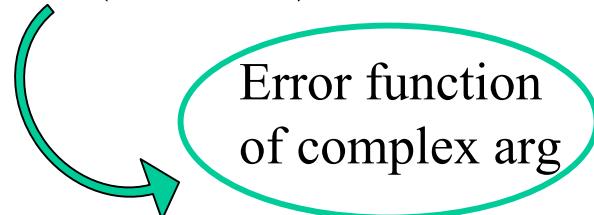
$$\partial_t f = p \partial_q f \quad \Rightarrow \quad \partial_t (\partial_p f) = \partial_p (p \partial_q f)$$

$$\partial_p f(t + \Delta t) = \partial_p f(t) + \Delta t \left(\frac{(p \partial_q f)_{j+1} - (p \partial_q f)_{j-1}}{2\Delta p} \right)$$

In Boussard We Trust.

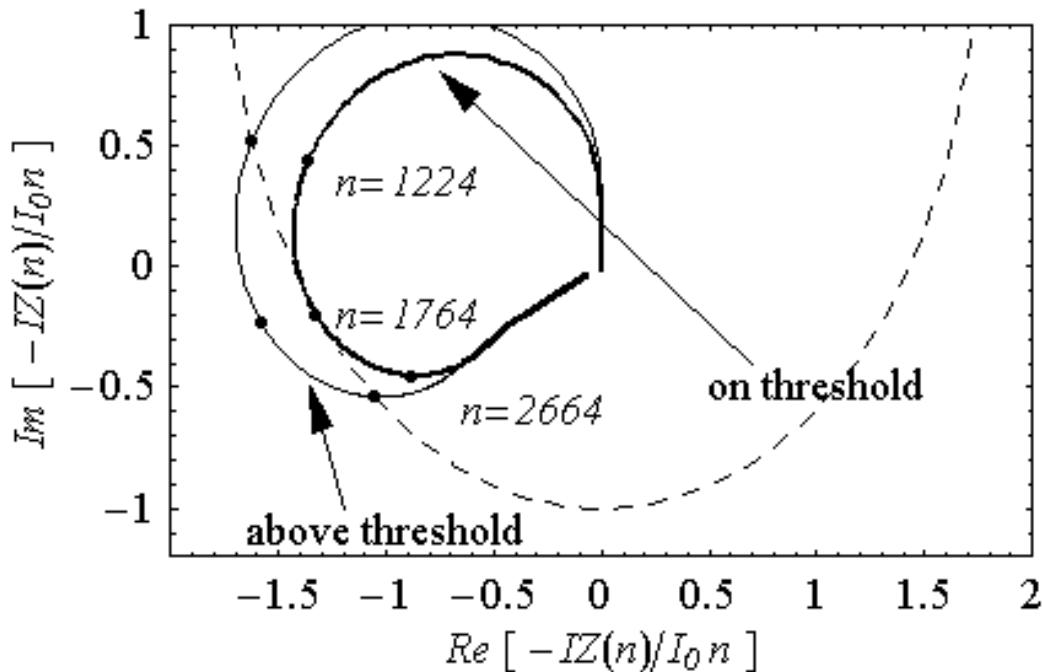
Stability Analysis of Linear Motion

- Boussard criterion (use equivalent coasting beam with same peak current)
- Gaussian distribution in energy spread
- Dispersion relation $\frac{I_c \omega_0 R^2}{\sigma_z^2 \sqrt{2\pi}} \frac{Z(n)}{n} = \frac{i}{W(vR / \sigma_z n)}$



Keil-Schnell Diagram

NSLS-VUV



Most unstable harmonic:
 $n = 1764 \Leftrightarrow \lambda = 6.7 \text{ mm}$

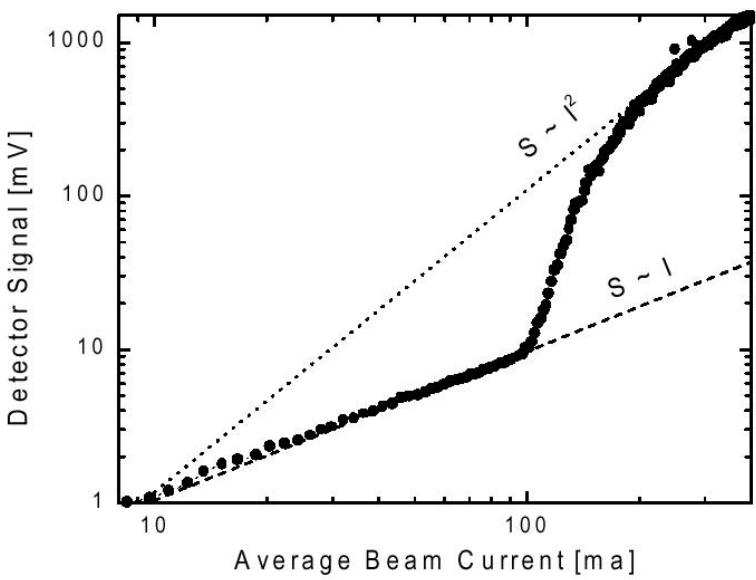
	Measurements *	Boussard	Solver
Current threshold	100 mA $I_c = 3.7 \text{ pC/V}$ $N = 1.1 \times 10^{11}$	$I_c = 6.2 \text{ pC/V}$ $N = 1.8 \times 10^{11}$	$I_c = 6.6 \text{ pC/V}$ $N = 1.9 \times 10^{11}$
CSR wavelength	7 mm	6.7 mm	7.8 mm

* G. Carr *et al.*, NIM-A 463 (2001) p. 387

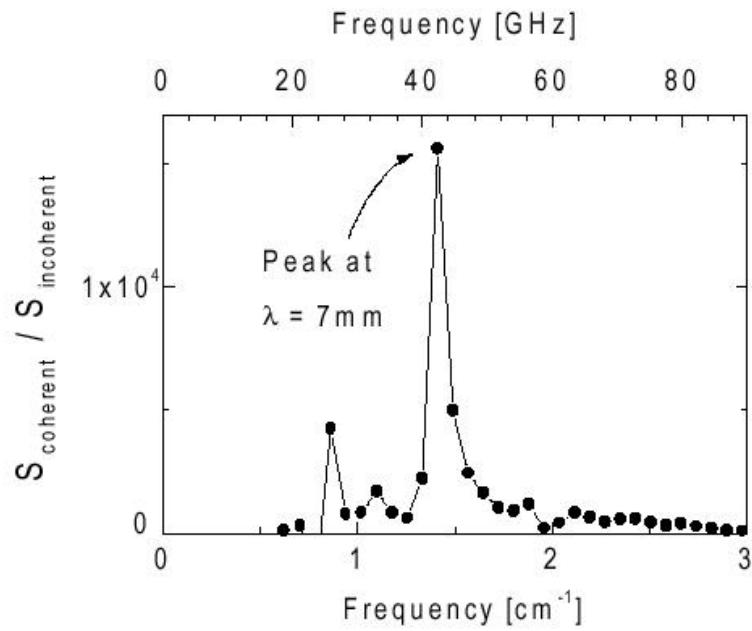
Observations of CSR - NSLS VUV Ring

Carr *et al.* NIM-A 463 (2001) p. 387

Current Threshold for Detection
of Coherent Signal

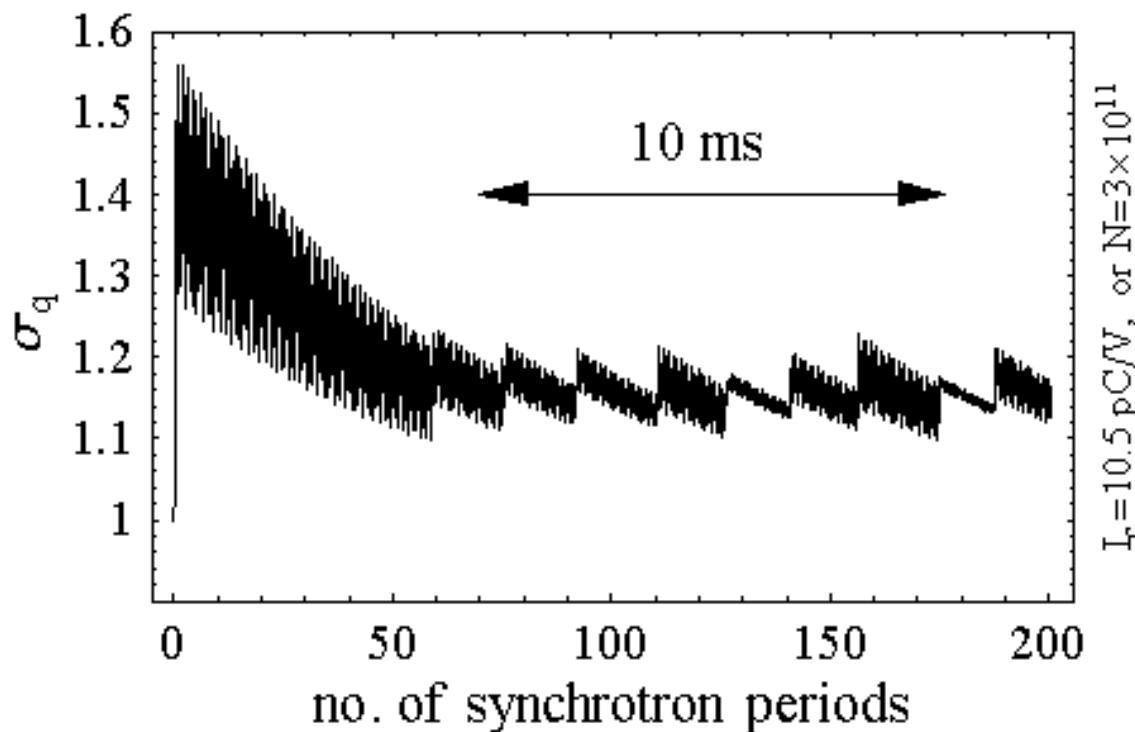


Spectrum of CSR Signal
(wavelength ~ 7 mm)



It Bursts!

Bunch Length

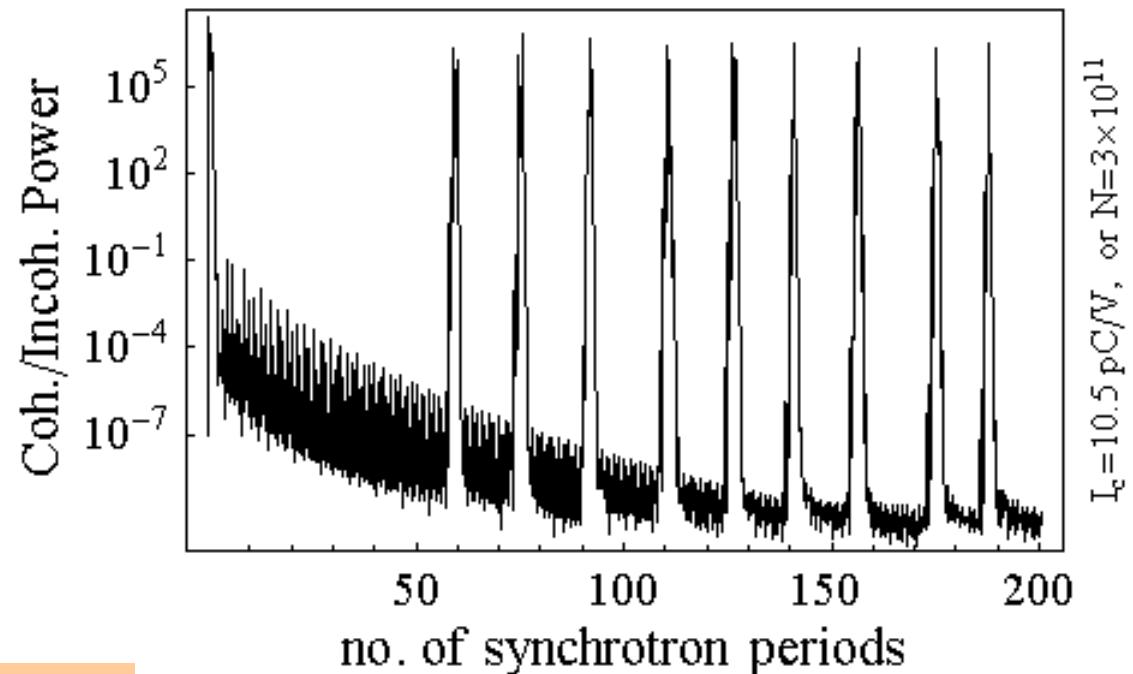


$$N = 3 \times 10^{11} \equiv I_c = 10.5 \text{ pC/V}$$

$$P_n^{coh} = 2(eN\omega_0)^2 \operatorname{Re} Z(n) |\lambda_n|^2$$

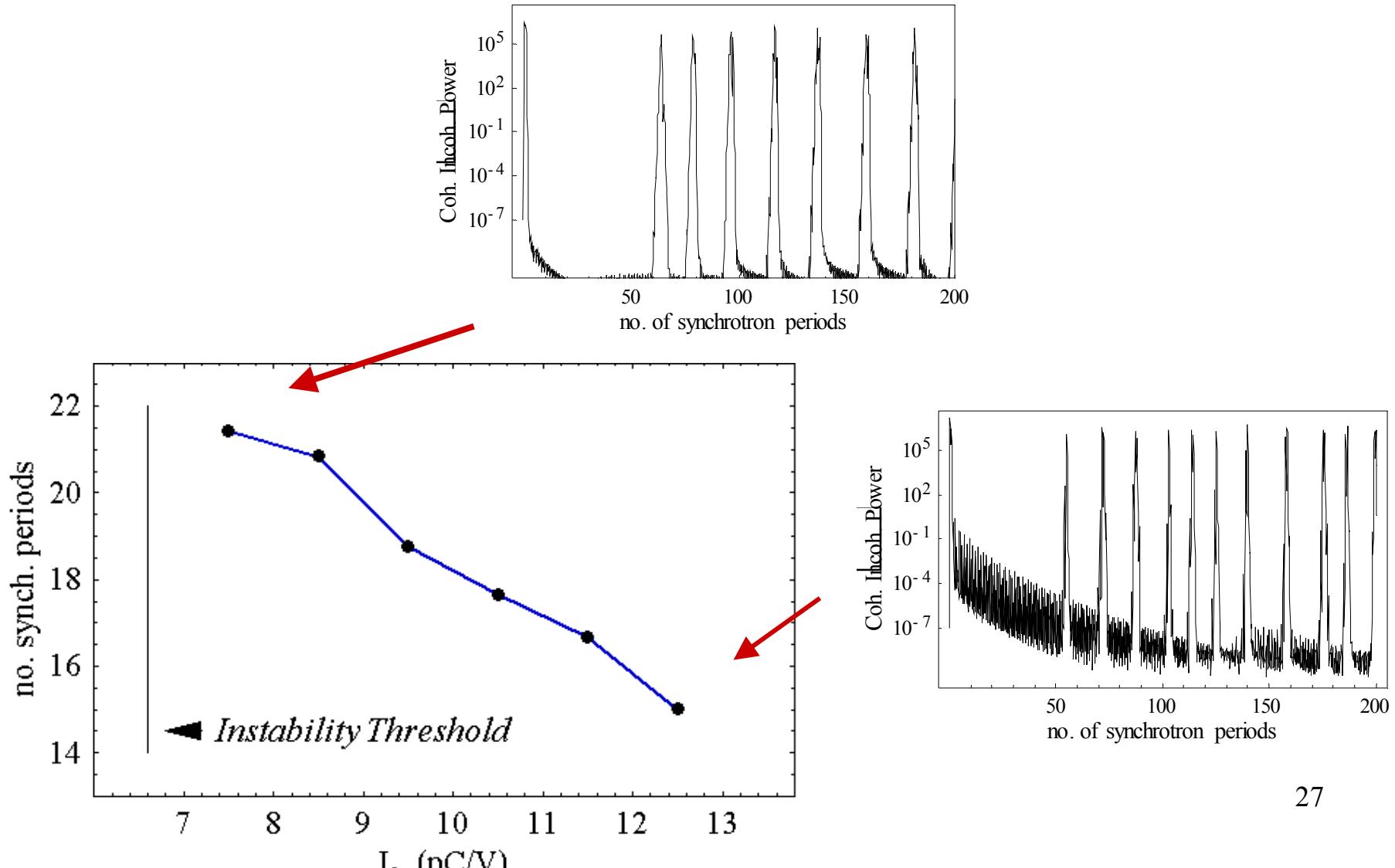
$$P_n^{incoh} = 2N(e\omega_0)^2 \operatorname{Re} Z(n) / (2\pi)^2$$

Coherent/Incoherent Power



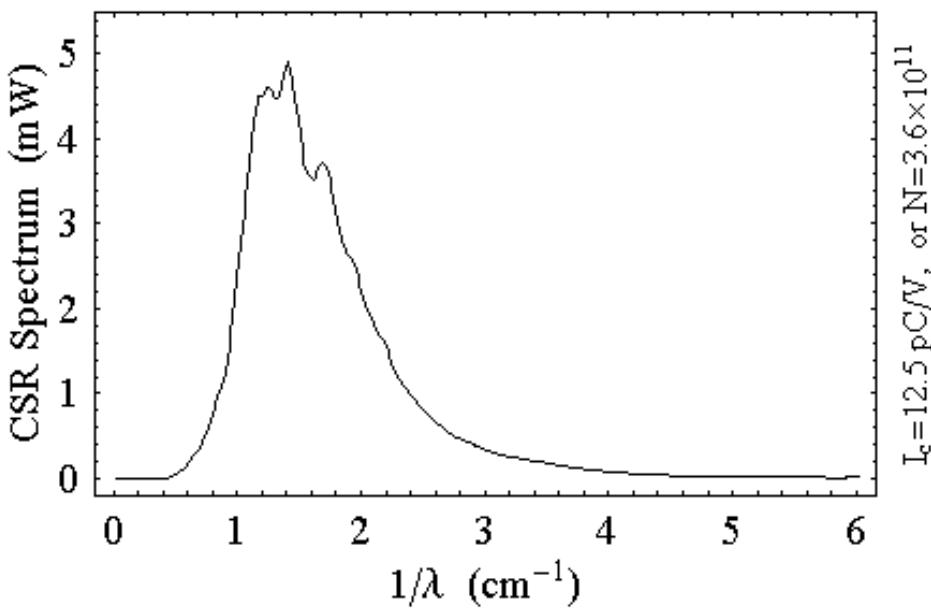
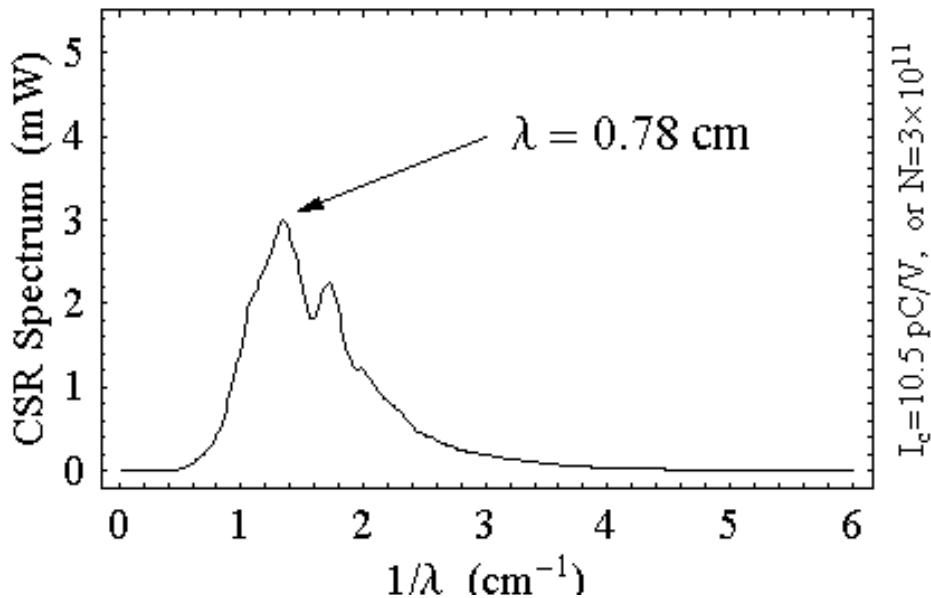
$$N = 3 \times 10^{11} \equiv I_c = 10.5 \text{ } pC/V$$

Bursts Separation vs. Current



Averaged CSR Emission

NSLS-VUV

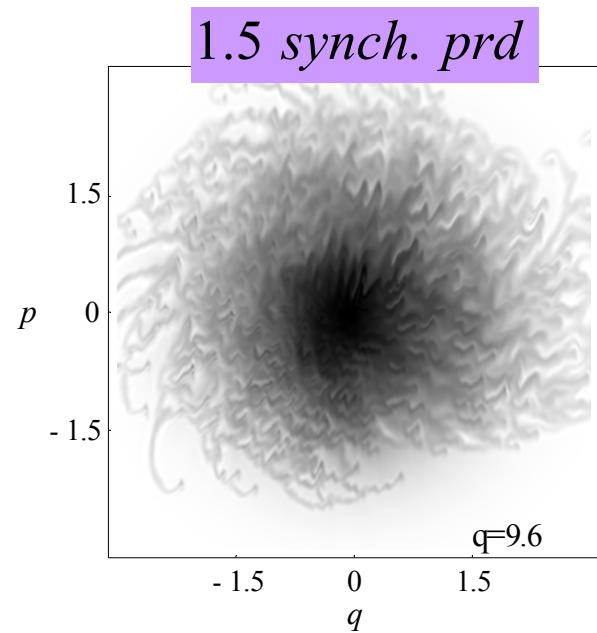
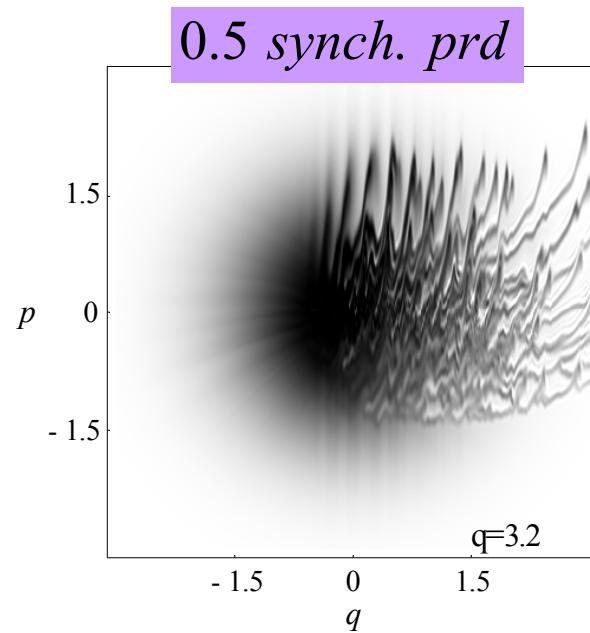
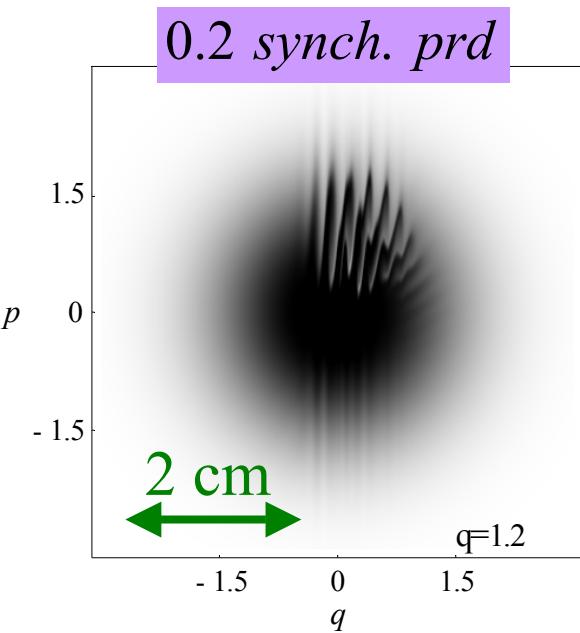


$$N = 3 \times 10^{11}$$

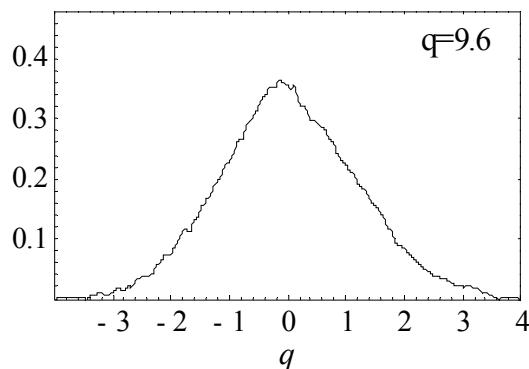
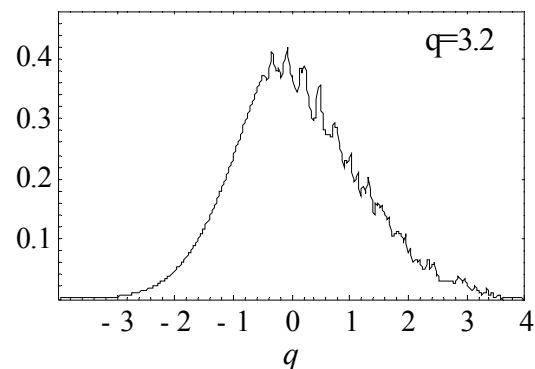
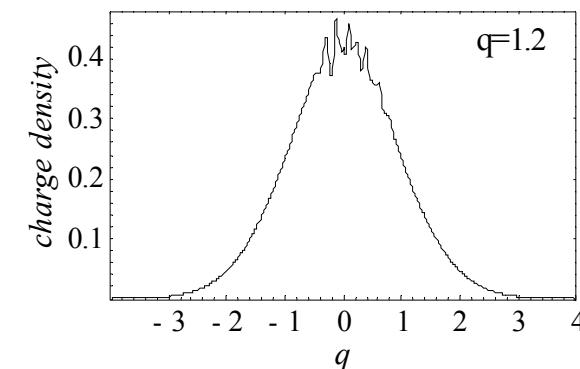
$$N = 3.6 \times 10^{11}$$

No damping, 25% > threshold

Density Plots in Phase space



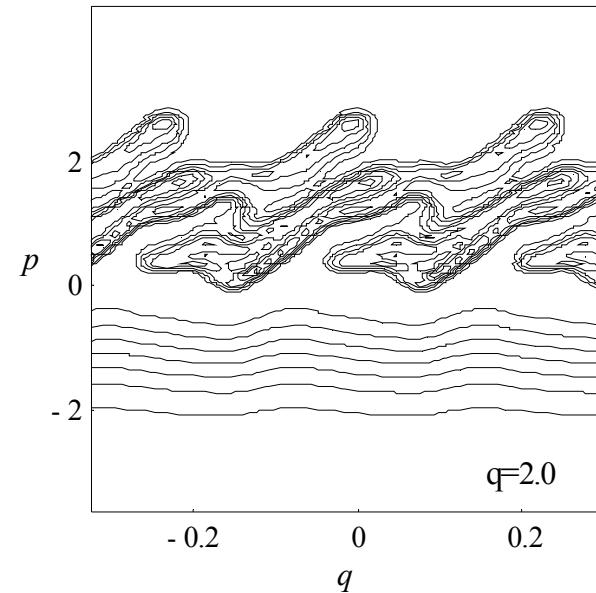
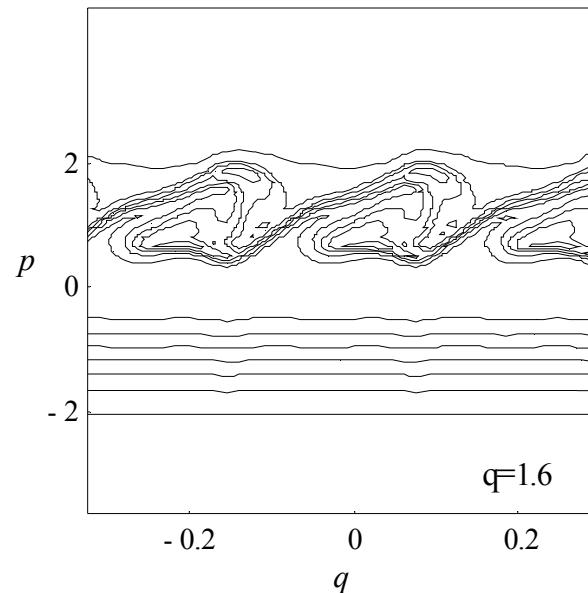
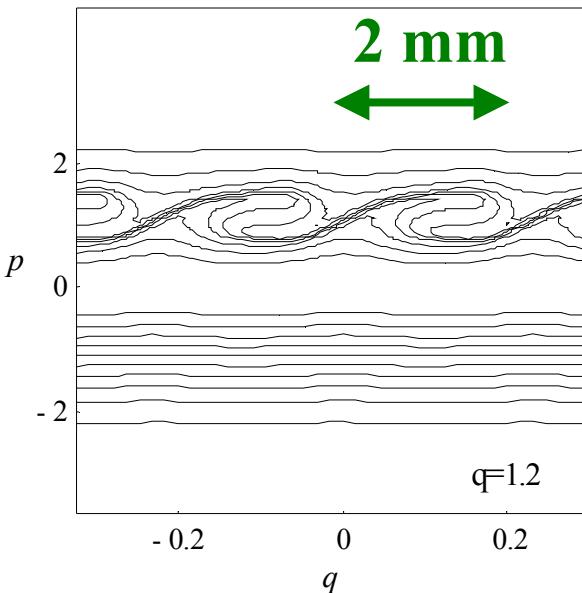
Charge Distribution



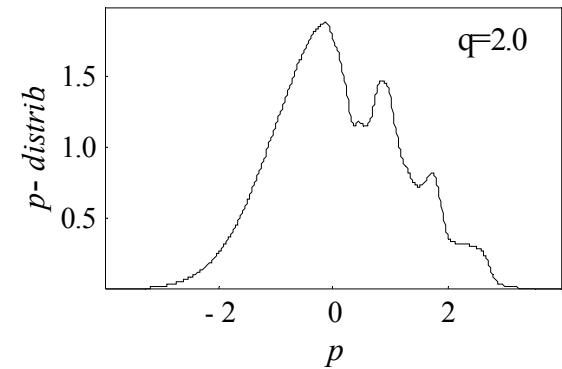
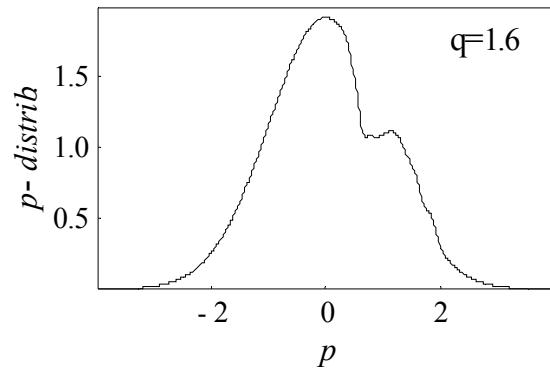
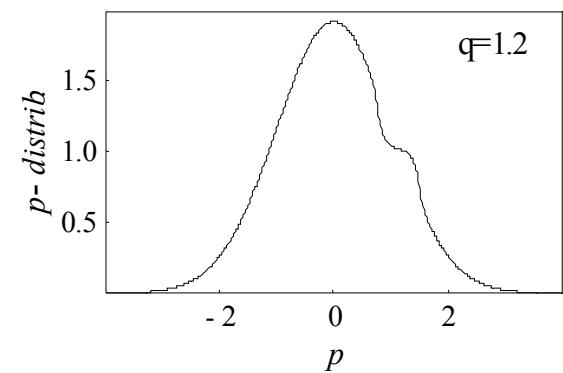
Coasting Beam

No damping, 25% > threshold

Density Contours in Phase space

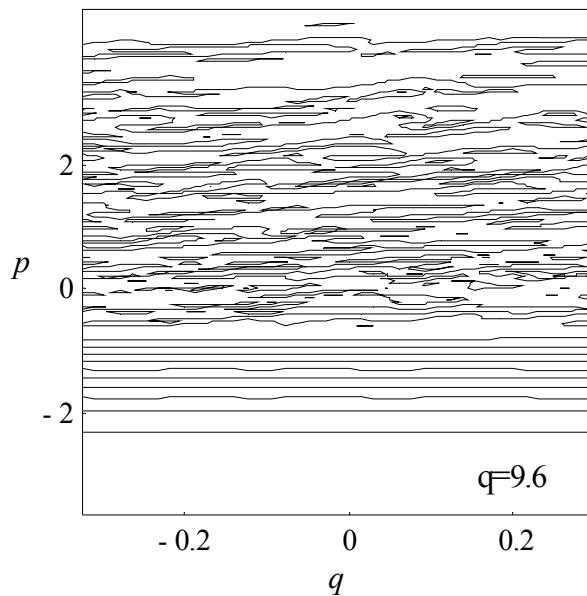


Energy Spread Distribution

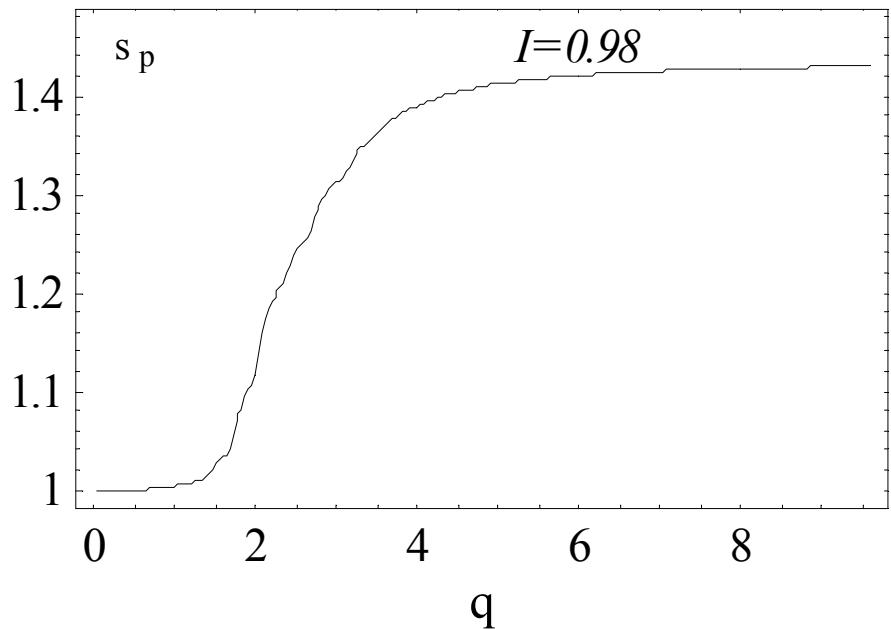


Coasting Beam: Asymptotic Solution

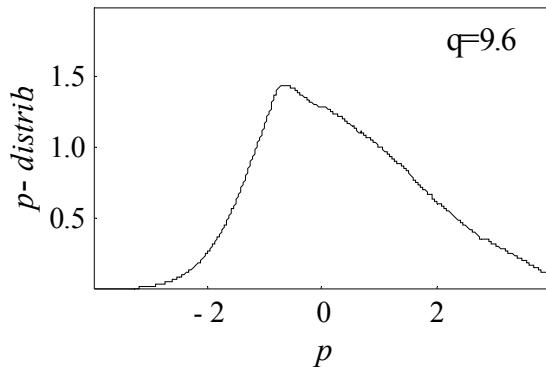
Density Contours in Phase space



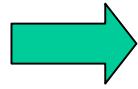
Energy Spread vs. Time



Energy Spread Distribution



Large scale structures have disappeared.



Distribution approaches some kind of steady state.

Conclusions

- Numerical model gives results **consistent with linear theory**, when this applies.
- CSR **instability saturates quickly**.
- **Saturation removes microbunching**, enlarges bunch distribution in phase space.
- Relaxation due to **radiation damping gradually restores conditions for CSR instability**.
- CSR instability + radiation damping = **sawtooth-like pattern & bursting** (qualitative consistency with observations).