

Details and Results of Additional Experiments

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I. Introduction

II. Streak Camera:

- Observed Bunch length $> 3\text{ps}$
- Asymmetric Bunch Shapes
- Explanation of Resolution Limit

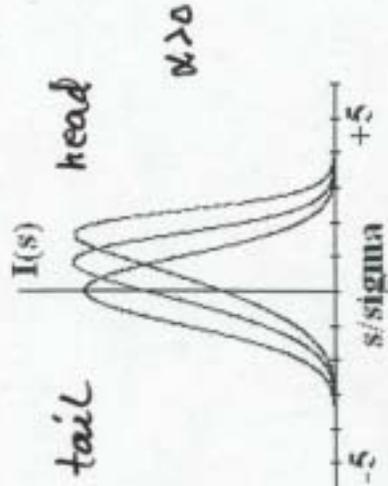
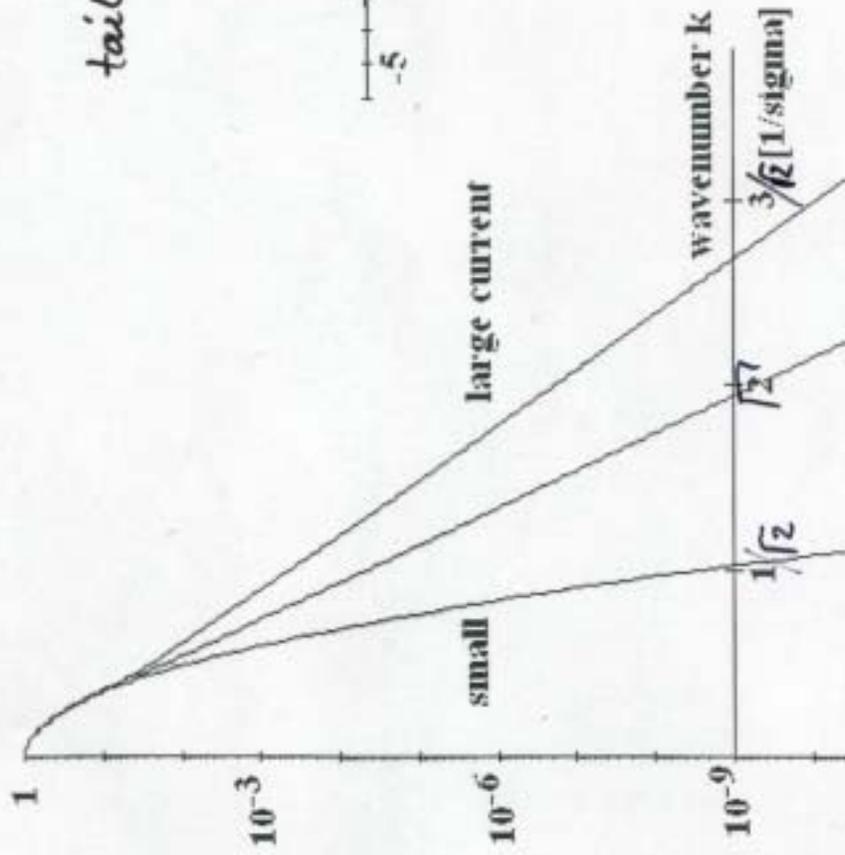
III. Absolute Measurement of the Energy Spread by Compton Backscattering: Microwave Threshold

IV. Determination of Higher Order Contributions to the Momentum Compaction Factor: i) From Measurements of the Synchrotron Tune and ii) The Orbits as a Function of the Pathlength

V. Longitudinal Beam Spectra and Bursting IR Signals CSR and/or Chamber Impedance?

VI. Conclusion

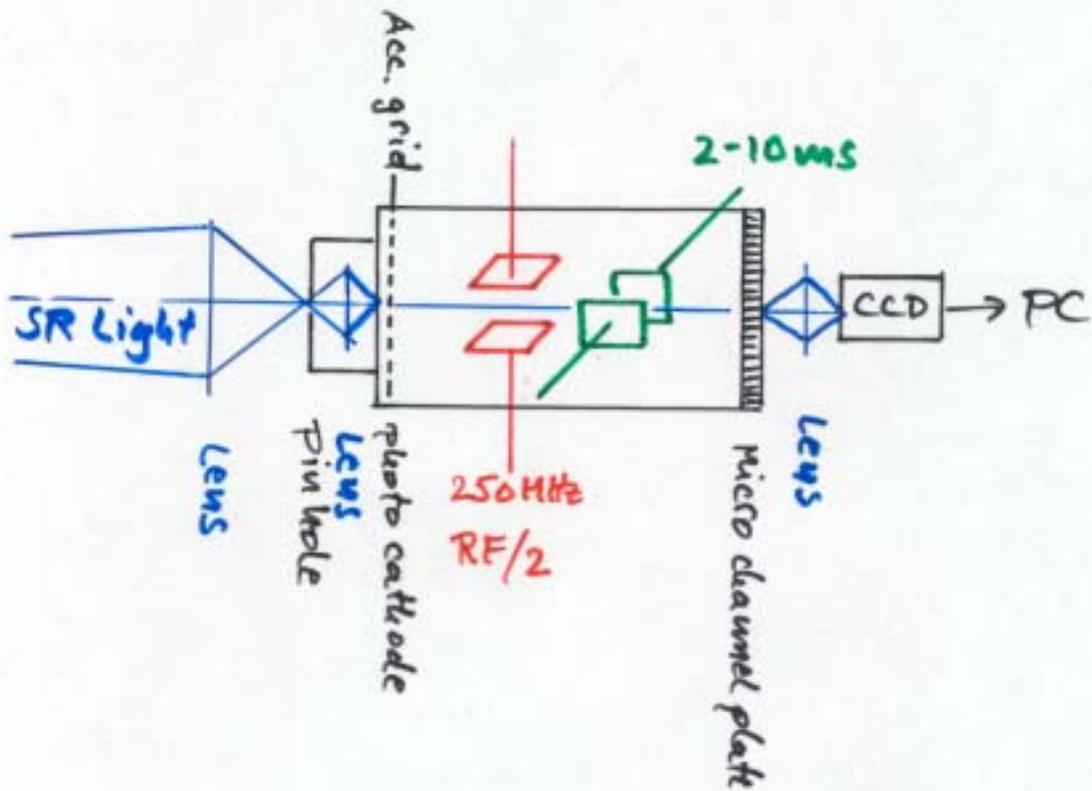
Formfactor with Potential Well Distortion by CSR and Resistive Impedance



(Tane, Krinsky, Murphy in AP367)

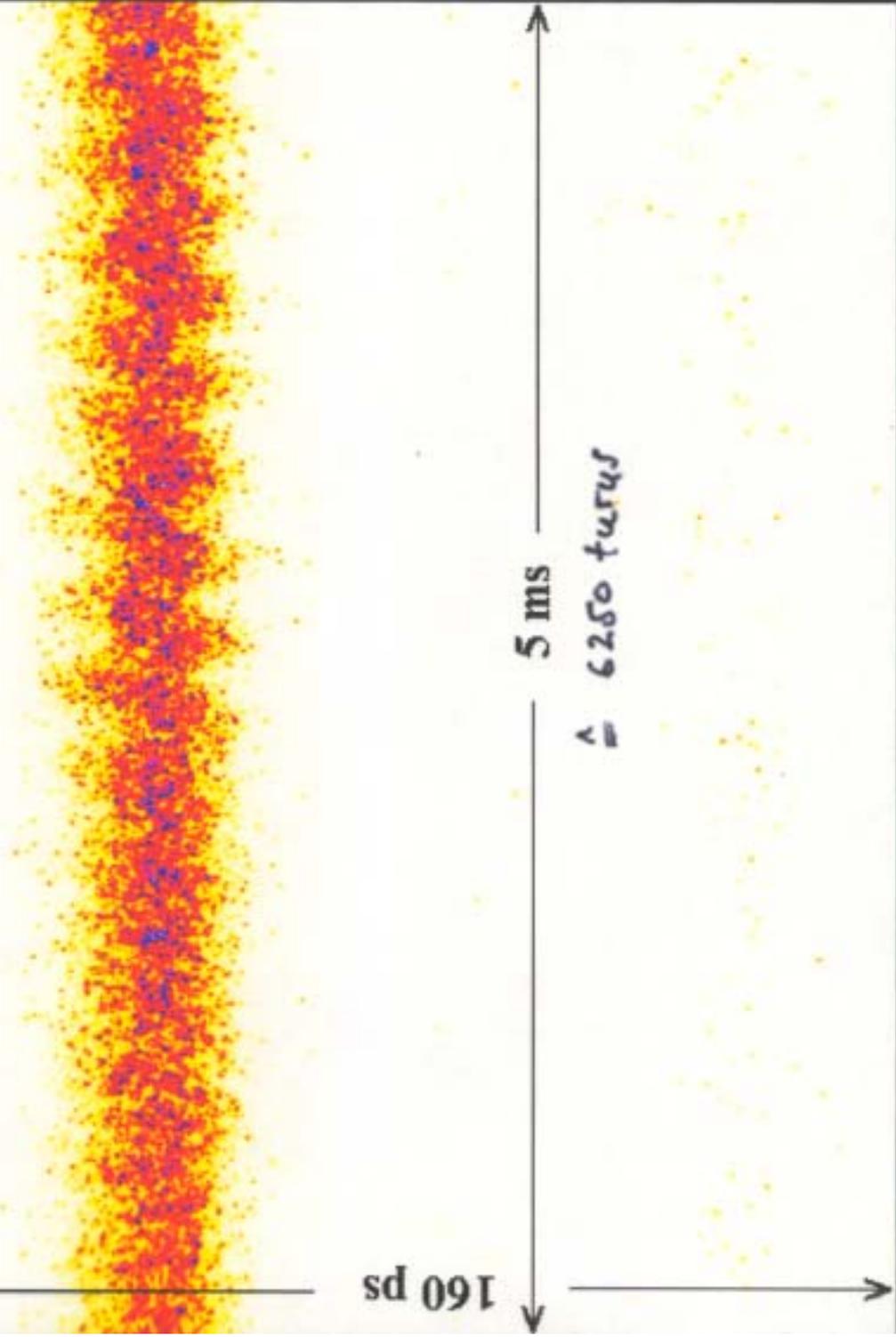
Streak Camera

Dual Sweep SC type C5680 (Hamamatsu)

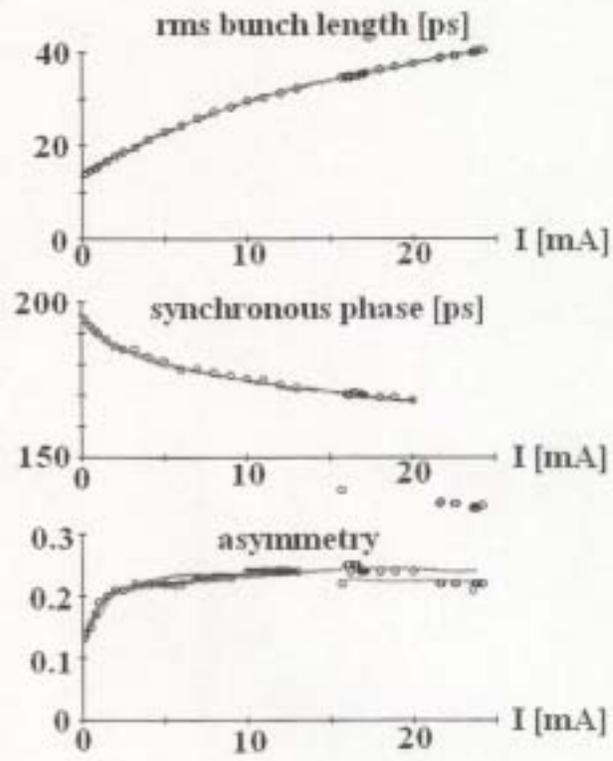


D:\D ateien\020922\1a\sb3kHz\0p056ma.img [Zoom x 1]

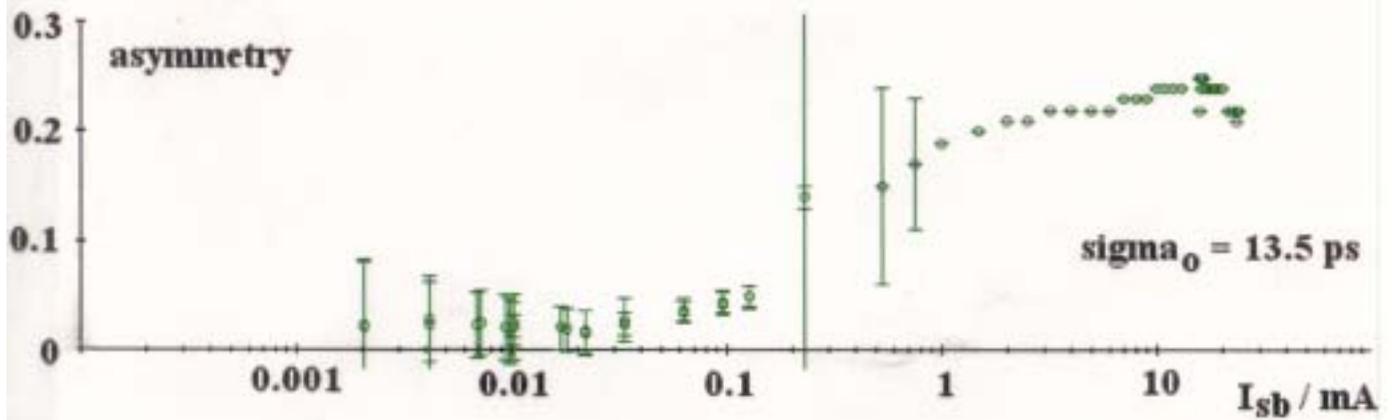
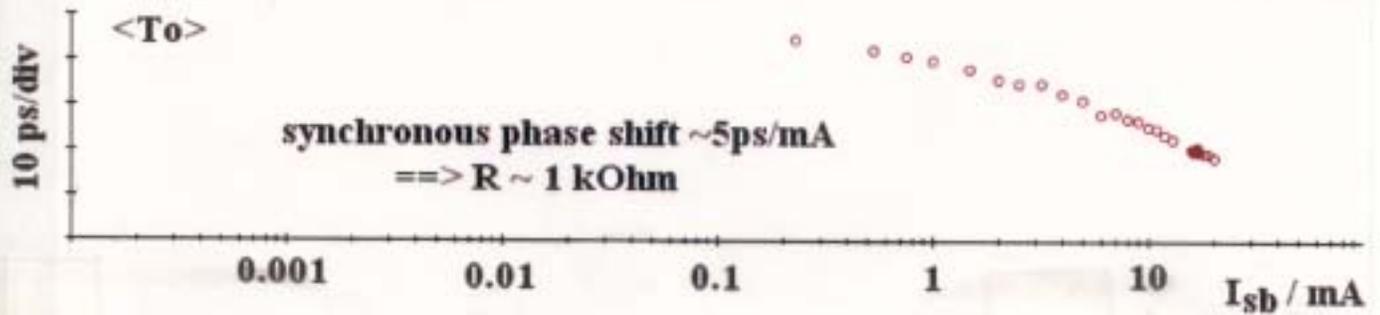
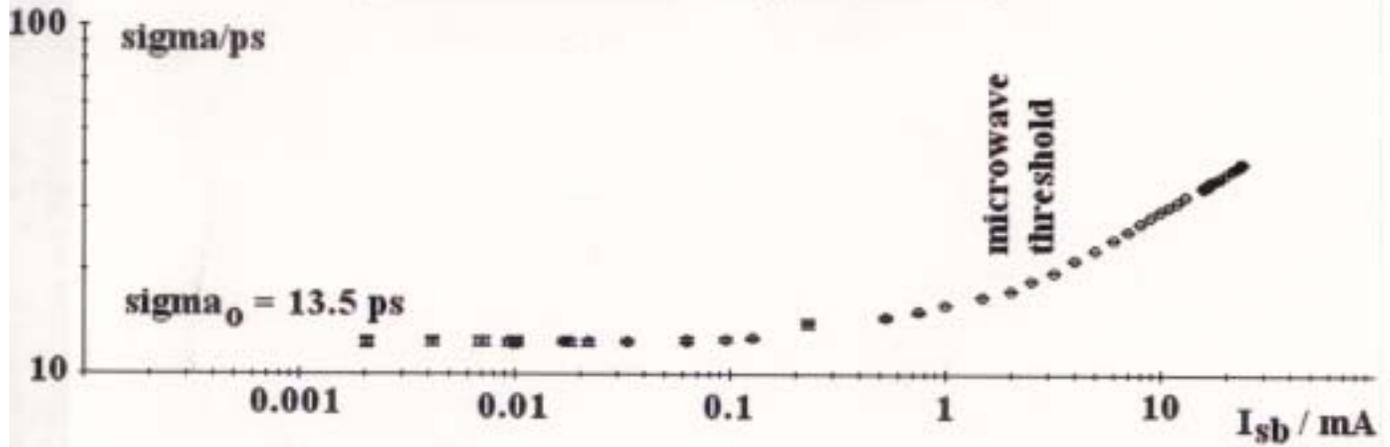
$I = 0.056 \text{ mA}$ $F_{\text{syn}} = 3 \text{ kHz}$



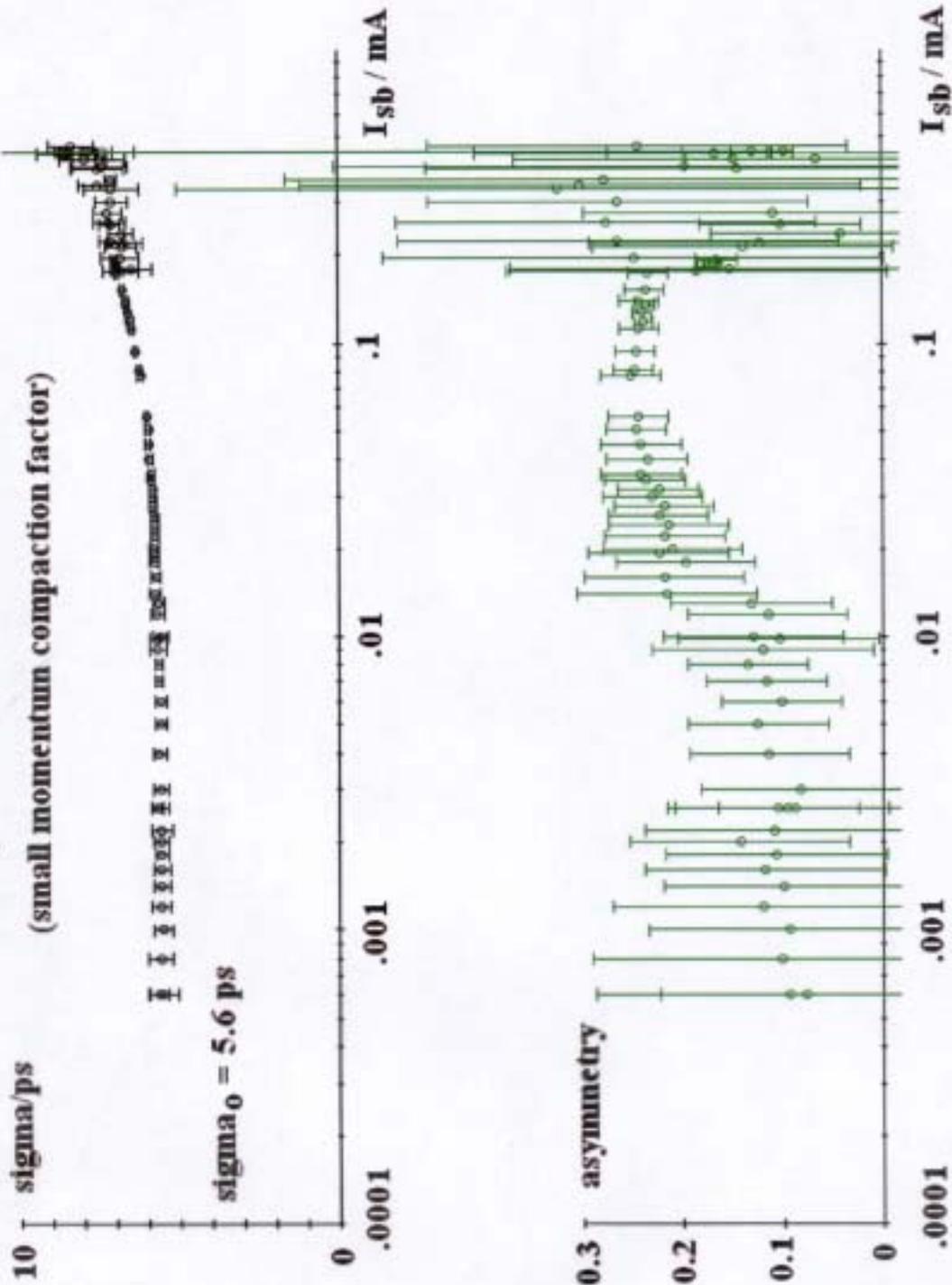
$$\alpha = 7.3 \cdot 10^{-4}$$



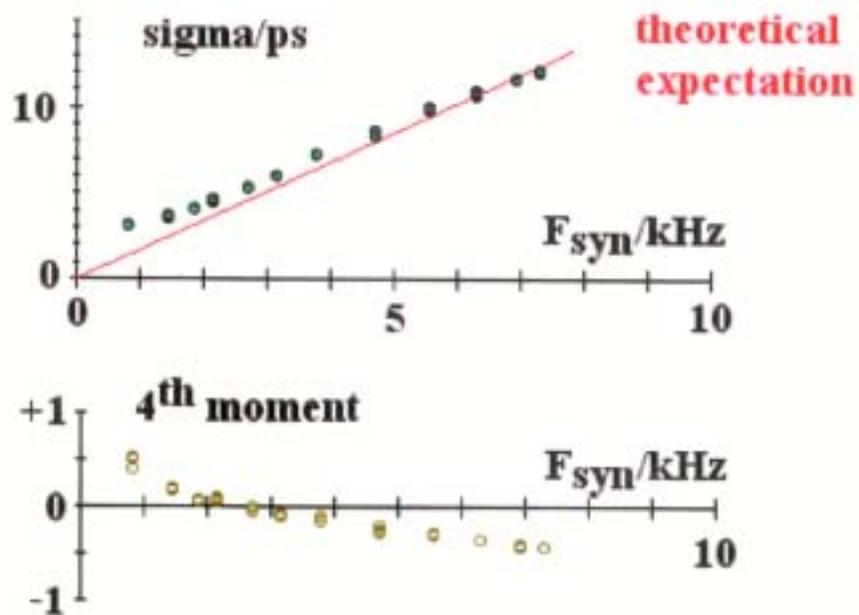
Analysis of Streak Camera Data (normal momentum compaction factor)



Analysis of Streak Camera Data (small momentum compaction factor)



**Bunchlength and 4th Moments of the Bunch Shape
as a Function of the Synchrotron Frequency
(with less than $2 \cdot 10^6$ electrons per bunch)**

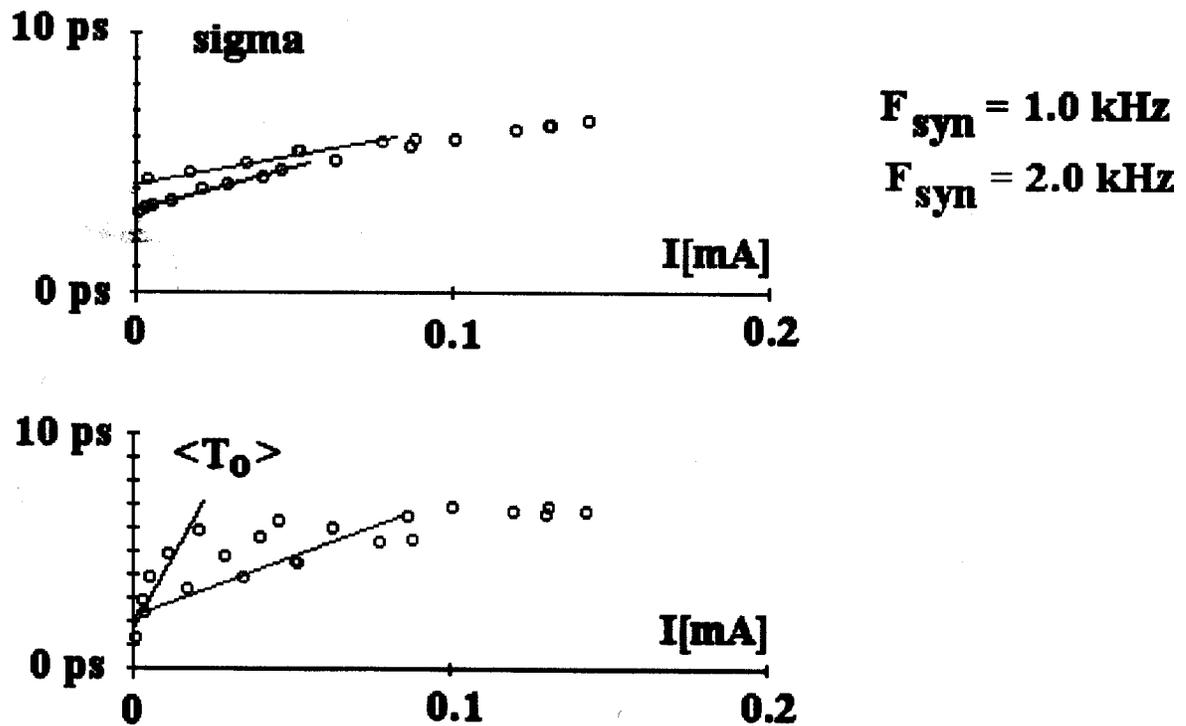


Conclusion and Question:

No bunch length shorter than 3 ps observed.

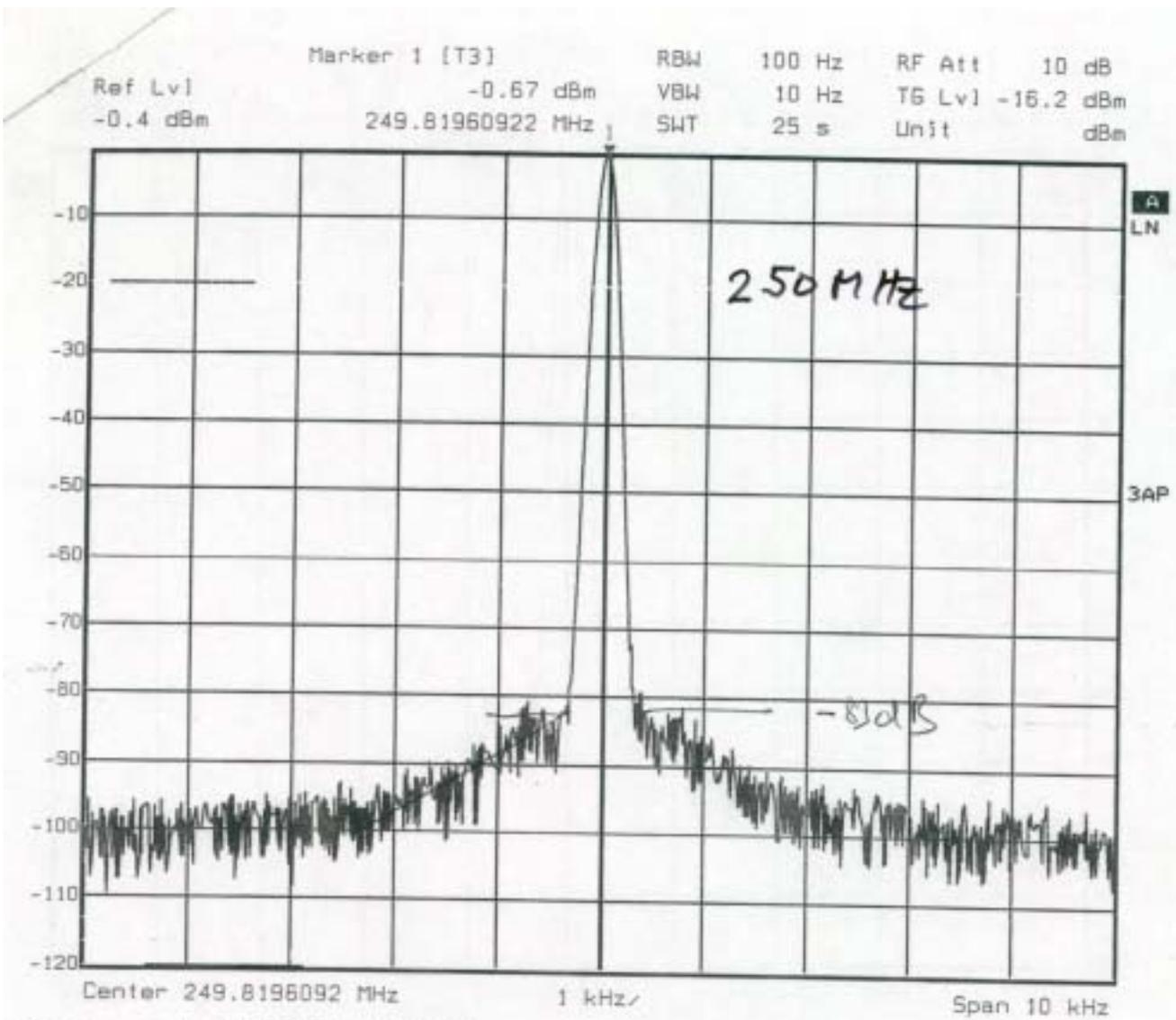
Resolution limit of the streak camera?

Bunchlength and Synchronous Phase as a Function of Beam Current



Conclusion:

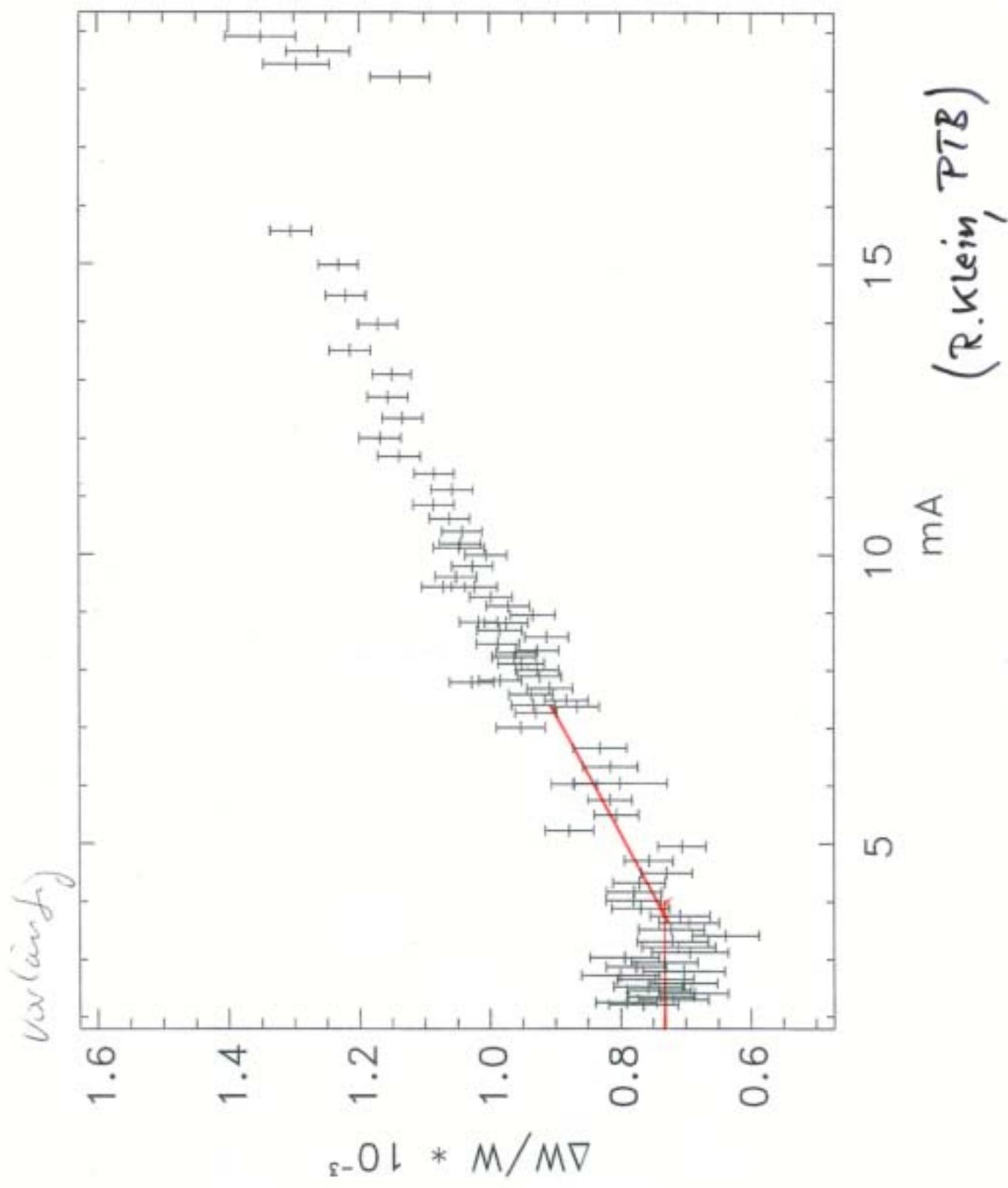
The synchronous phase shifts indicate bunches shorter than the resolution limit of our streak camera.



Date: 25.SEP.2002 11:42:25

⇒ Resolution Limitation of SC
due to Phase Noise

Messung 4.7.02



Momentum Compaction Factor from Measurements of the:

- i) Synchrotron Tune and
- ii) Orbits

as a Function of the Pathlength.

$$\frac{\Delta L}{L} = - \frac{\Delta r f}{r f} = \alpha \cdot \frac{\Delta P}{P}$$

$$\alpha = \alpha_1 + \alpha_2 \frac{\Delta P}{P} + \alpha_3 \left(\frac{\Delta P}{P}\right)^2 \dots$$

Make inverted expansion from:

$$\frac{\Delta r f}{r f} = - \sum_i \alpha_i \left(\frac{\Delta P}{P}\right)^i \rightarrow \frac{\Delta P}{P} = \sum_i \beta_i \left(\frac{\Delta r f}{r f}\right)^i$$

With the first few terms:

$$\frac{\Delta P}{P} = -\frac{1}{\alpha_1} \frac{\Delta r f}{r f} - \frac{\alpha_2}{\alpha_1^2} \left(\frac{\Delta r f}{r f}\right)^2 - \frac{2\alpha_2^2 - \alpha_1 \alpha_3}{\alpha_1^3} \left(\frac{\Delta r f}{r f}\right)^3 - \frac{\alpha_1^2 \alpha_4 + 5\alpha_2^3 - 5\alpha_1 \alpha_2 \alpha_3}{\alpha_1^4} \left(\frac{\Delta r f}{r f}\right)^4 \dots$$

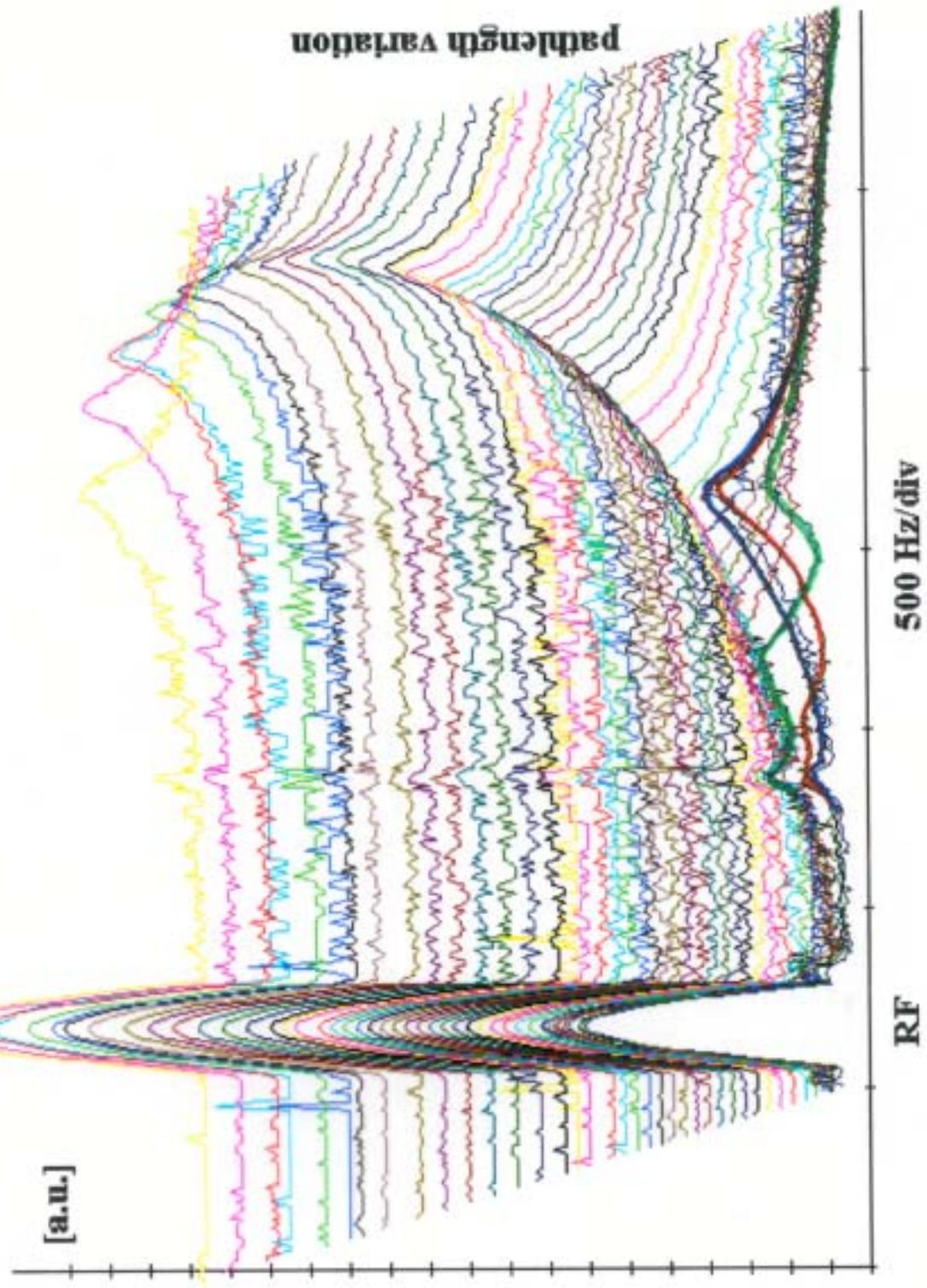
Synchrotron Tune depends on the local value of α :

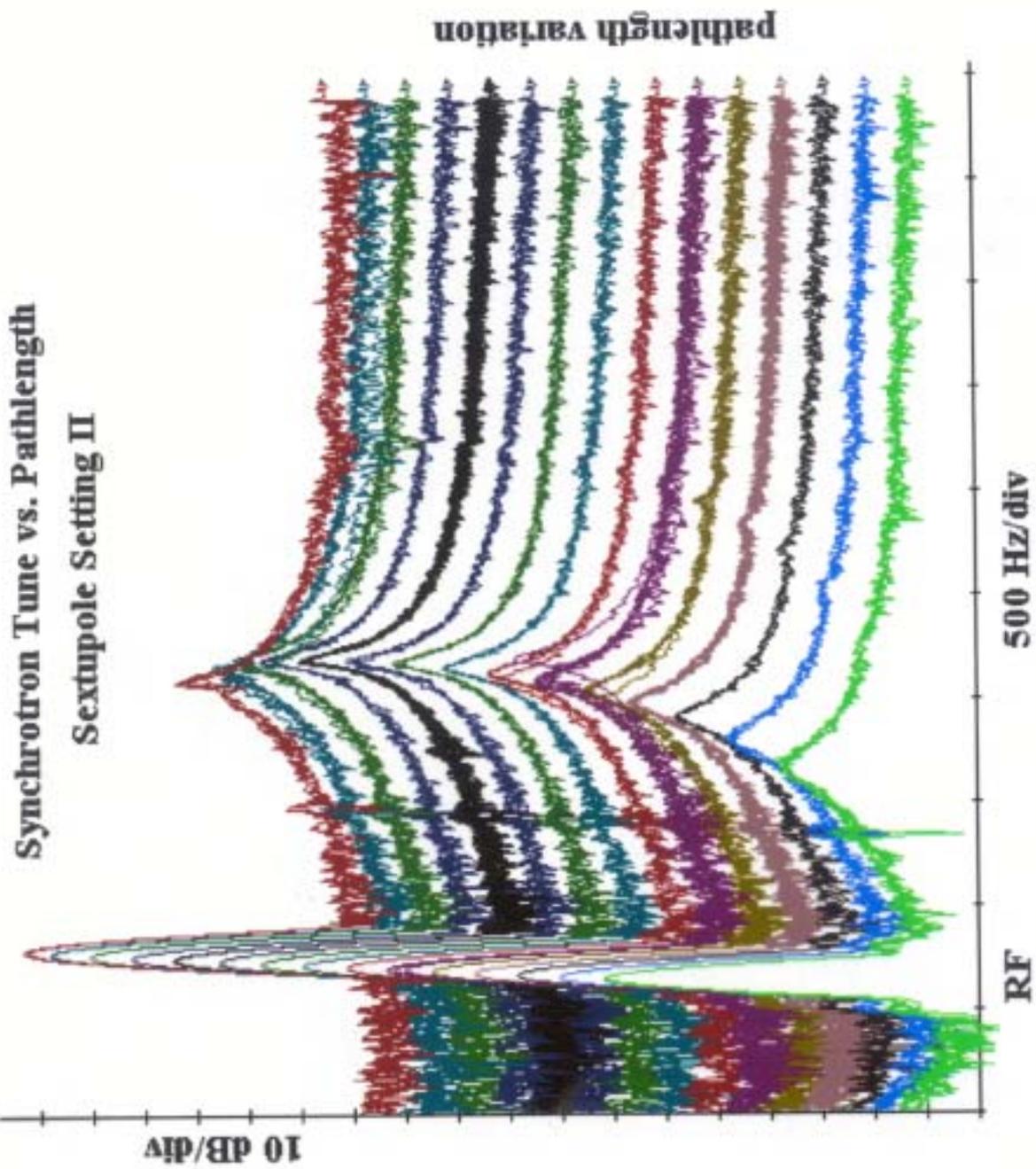
$$\alpha_{\text{local}} = \alpha_1 + 2\alpha_2 \frac{\Delta P}{P} + 3\alpha_3 \left(\frac{\Delta P}{P}\right)^2 + \dots$$

$$\Rightarrow \frac{\nu_{\text{syn}}^2}{\nu_{\text{syn}_0}^2} = 1 - \frac{2\alpha_2}{\alpha_1} \frac{\Delta r f}{r f} - \frac{2\alpha_2^2 - 3\alpha_1 \alpha_3}{\alpha_1^4} \left(\frac{\Delta r f}{r f}\right)^2 - \frac{4\alpha_2^3 - 8\alpha_1 \alpha_2 \alpha_3 + 4\alpha_1^2 \alpha_4}{\alpha_1^4} \left(\frac{\Delta r f}{r f}\right)^3 - \dots$$

(read: Frith...)

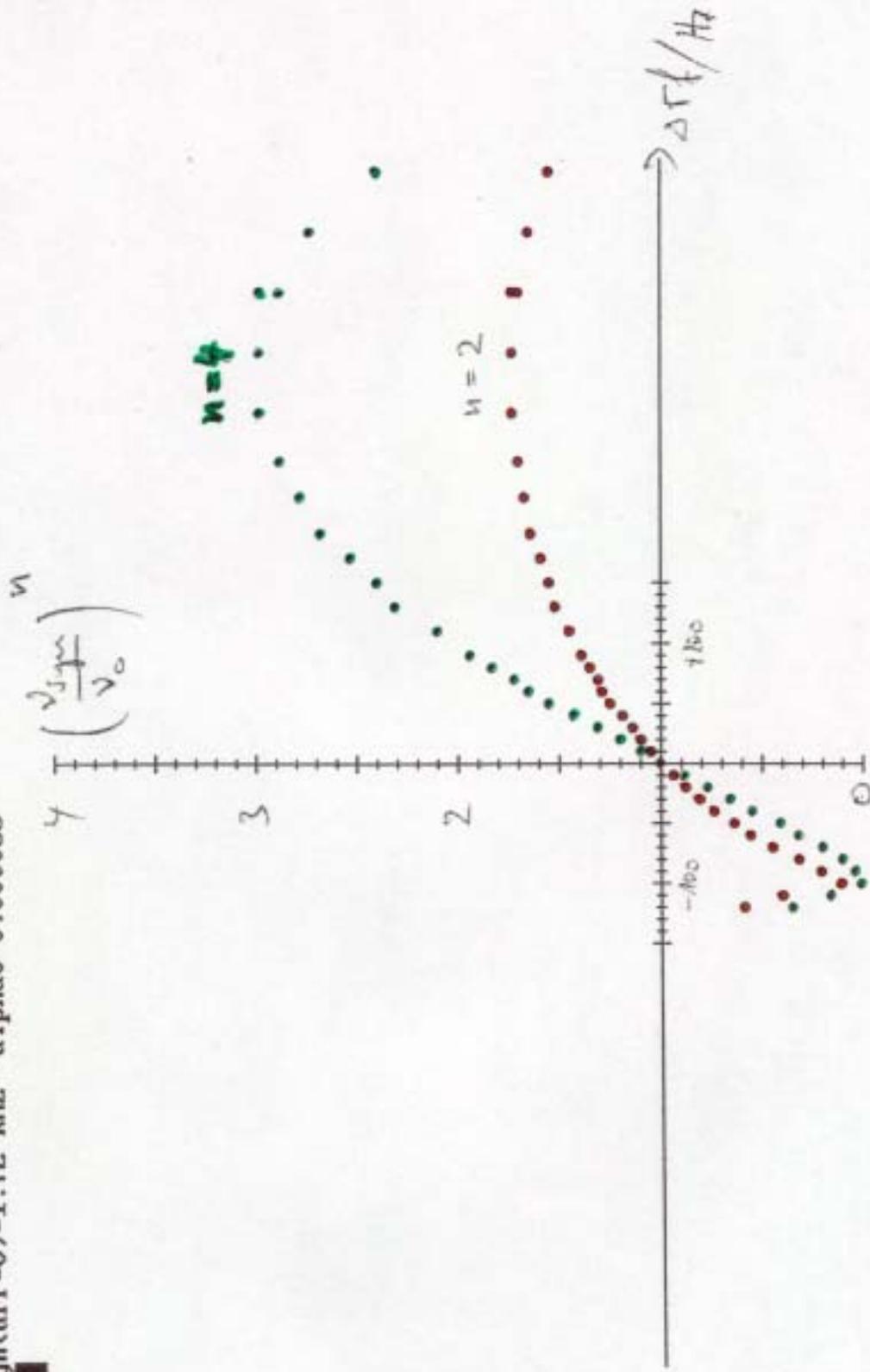
Synchrotron Tune vs. Pathlength Sextupole Setting I





Sextupole Setting I

a:ms020428.dat
Fsyn(drf=0)=1.72 kHz alpha0=0.000038
?



Orbit Distortions:

$$\Delta_{x,y} = D_{x,y} \cdot \frac{\Delta p}{p}$$

$$D = D_1 + D_2 \frac{\Delta p}{p} + D_3 \left(\frac{\Delta p}{p}\right)^2 + \dots$$

With

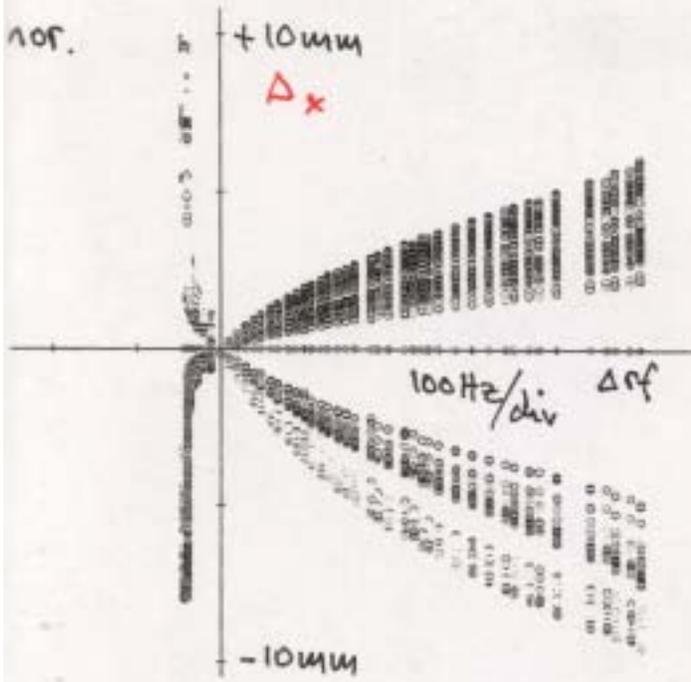
$$\frac{\Delta p}{p} = -\frac{1}{\alpha_1} \frac{\Delta r_f}{r_f} - \frac{\alpha_2}{\alpha_1^3} \left(\frac{\Delta r_f}{r_f}\right)^2 - \frac{2\alpha_2^2 - \alpha_1 \alpha_3}{\alpha_1^5} \left(\frac{\Delta r_f}{r_f}\right)^3 \dots$$

+ Neglecting D_3 and higher order contributions to the Dispersion one can write:

$$\alpha_1 \frac{\Delta}{D_1} + \alpha_2 \left(\frac{\Delta}{D_1}\right)^2 + \alpha_3 \left(\frac{\Delta}{D_1}\right)^3 + \dots \approx -\frac{\Delta r_f}{r_f} + \frac{D_2}{\alpha_1 D_1} \left(\frac{\Delta r_f}{r_f}\right)^2 \dots$$

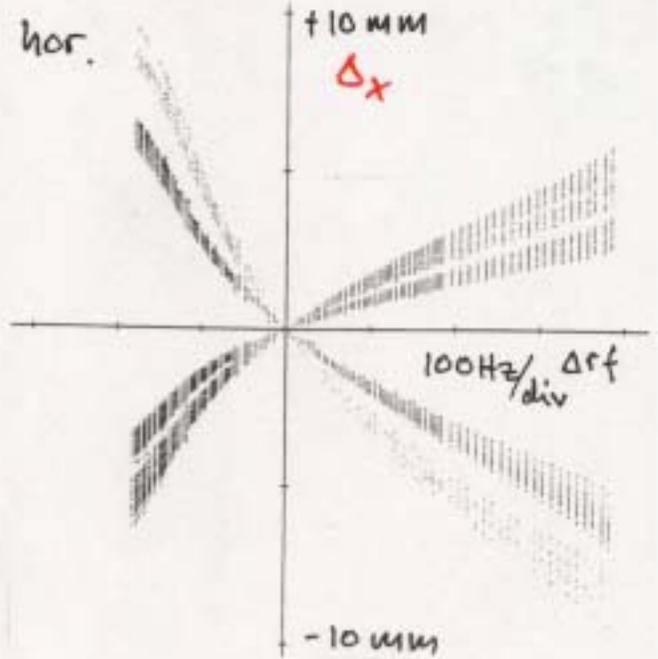
Off momentum Orbits $\Delta \left(\frac{\Delta r_f}{r_f} \right)$

hor.



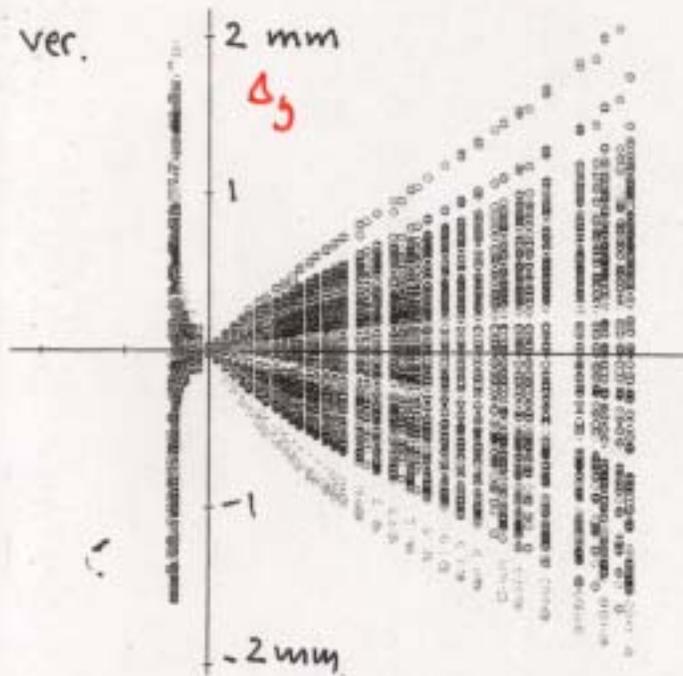
02/05/20

hor.

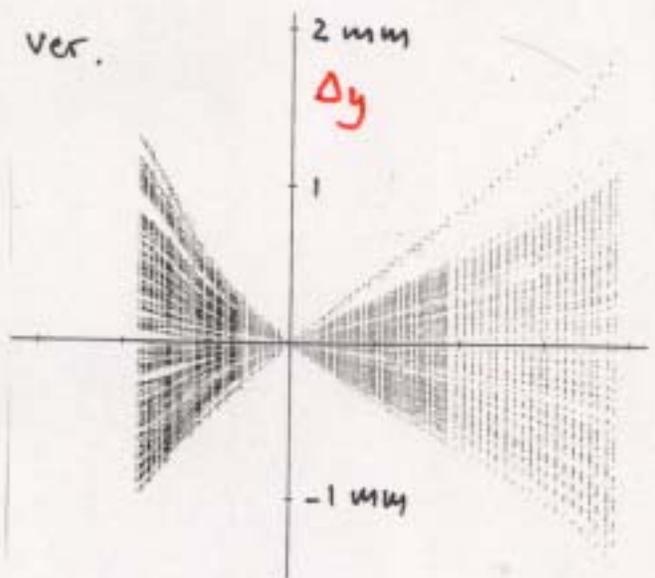


02/05/06

ver.



ver.



Fit results:

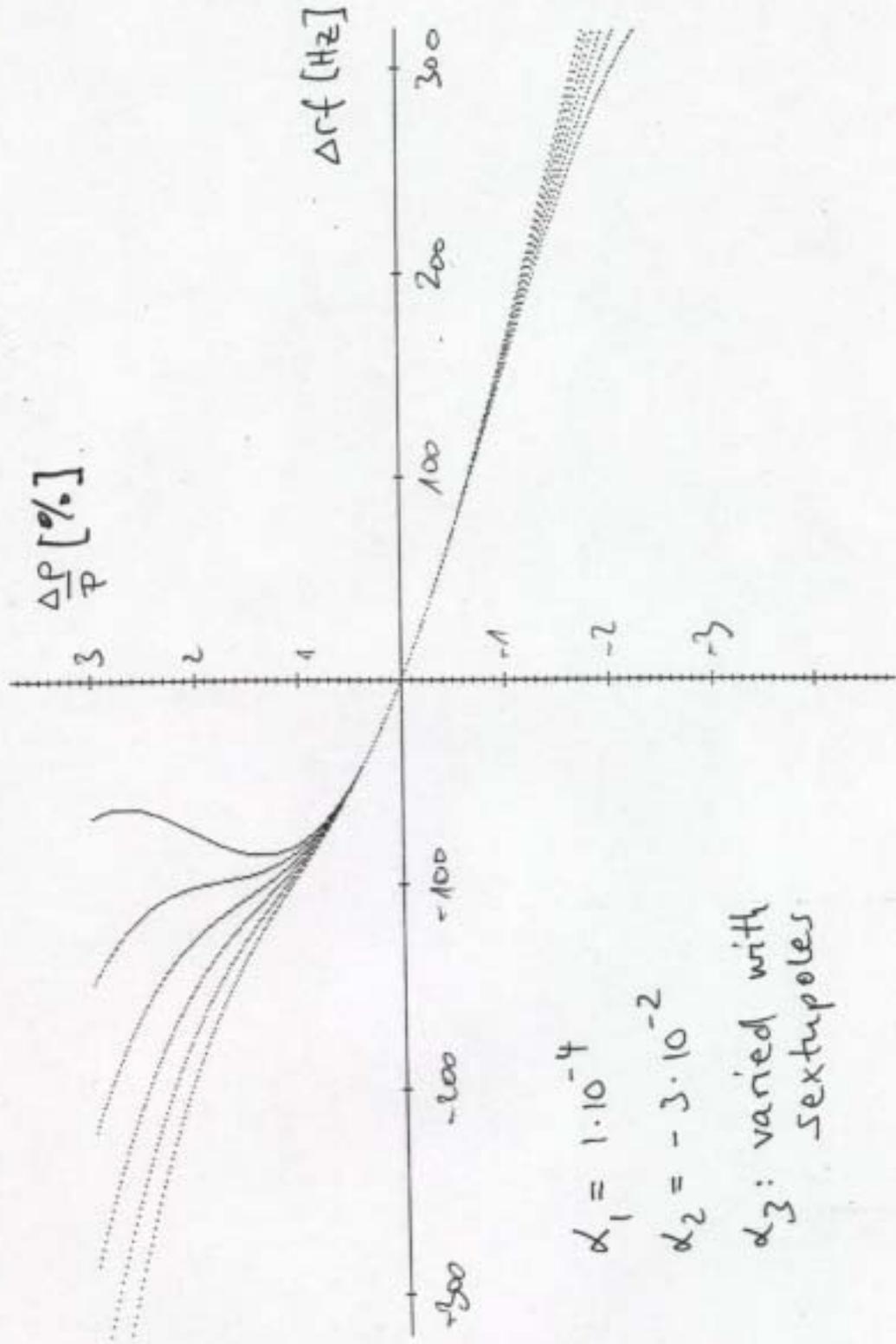
$$\alpha_1 = 3.5 \cdot 10^{-5}$$

$$\alpha_2 = -0.00110 (\pm 0.00017)$$

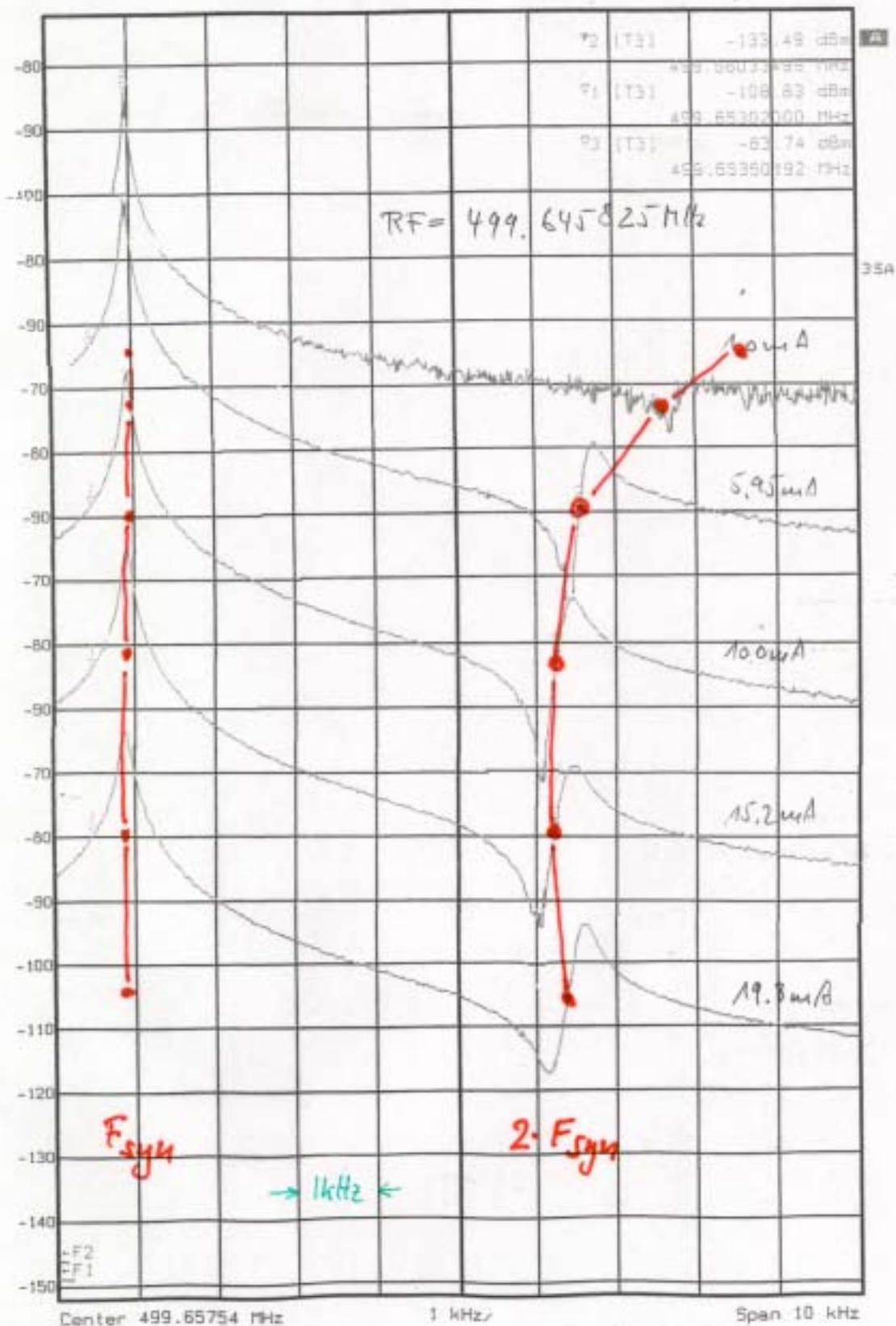
$$\alpha_3 = -0.024 (\pm 0.013)$$

$$\alpha_4 = 1.9 (\pm 0.5)$$

Results from Optics Calculations



IV. Measurement of Incoherent Sydnr. Tune shift with Intensity



Date: 5.NOV.2001 20:51:24

19.8 μ A

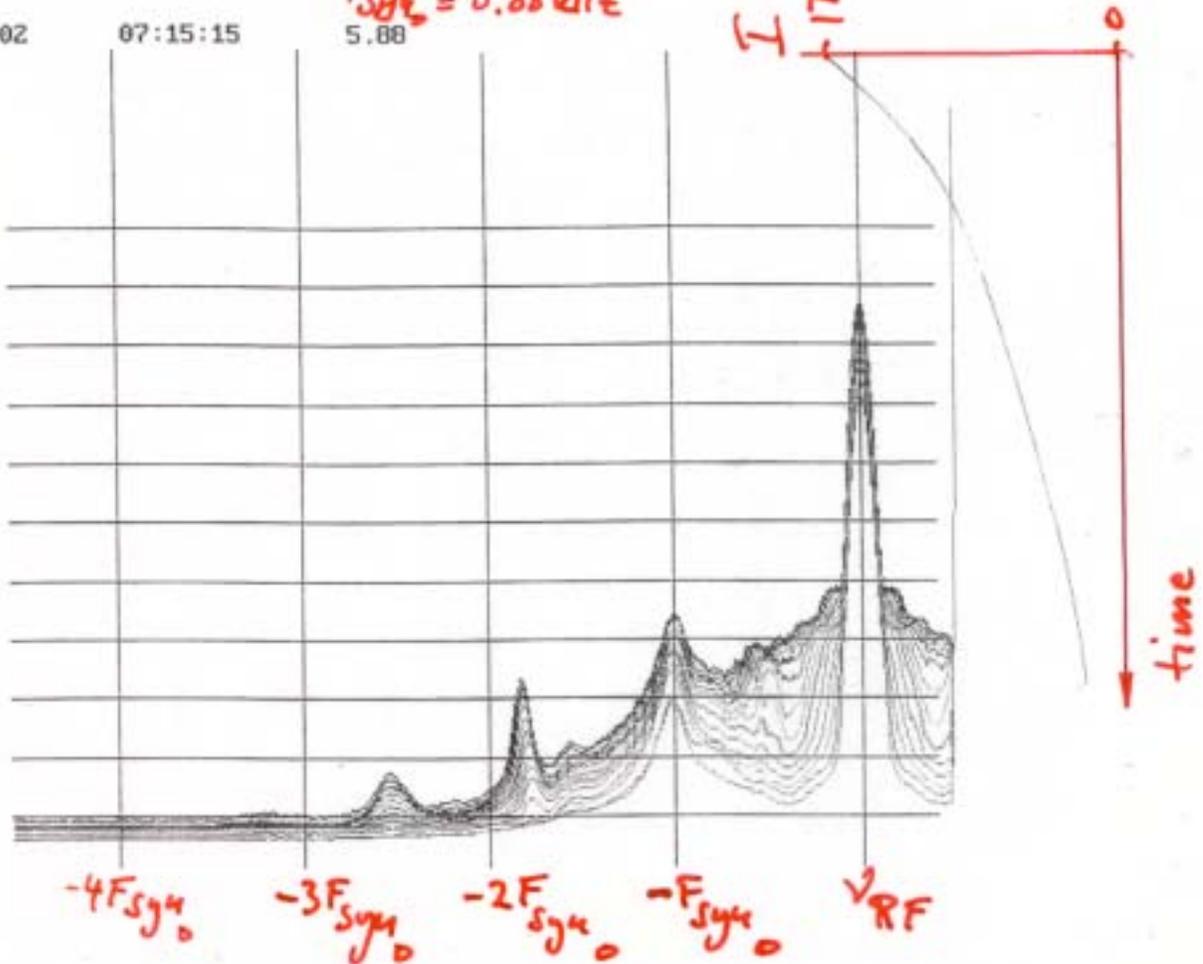
Longitudinal Modes with Single Bunch

07.07.02
?

07:15:15

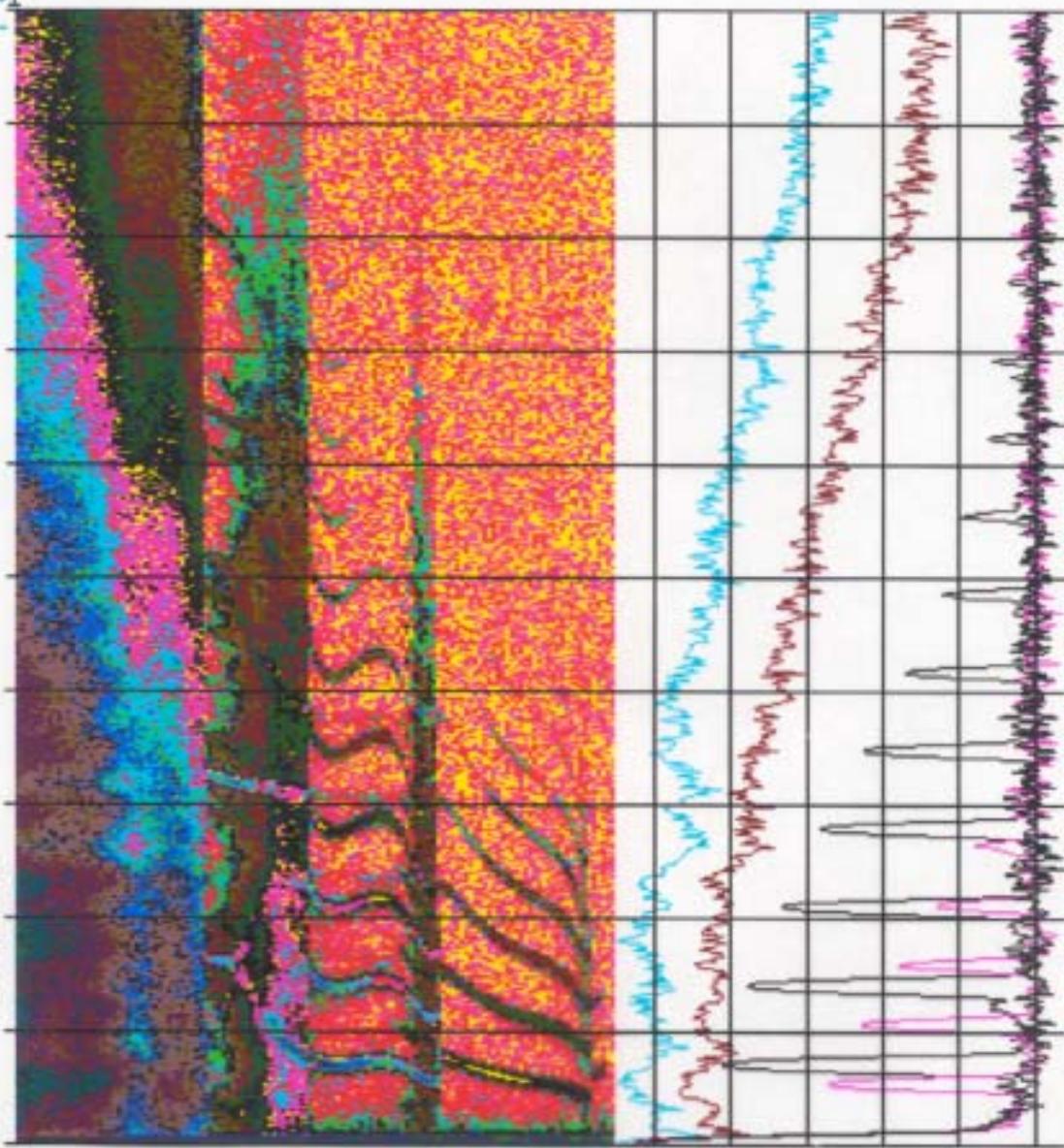
$F_{syn} = 5.88 \text{ kHz}$
5.88

$I = 17.5 \text{ mA}$



Bursting CSR - single Burst

12.09.02 03:26:01



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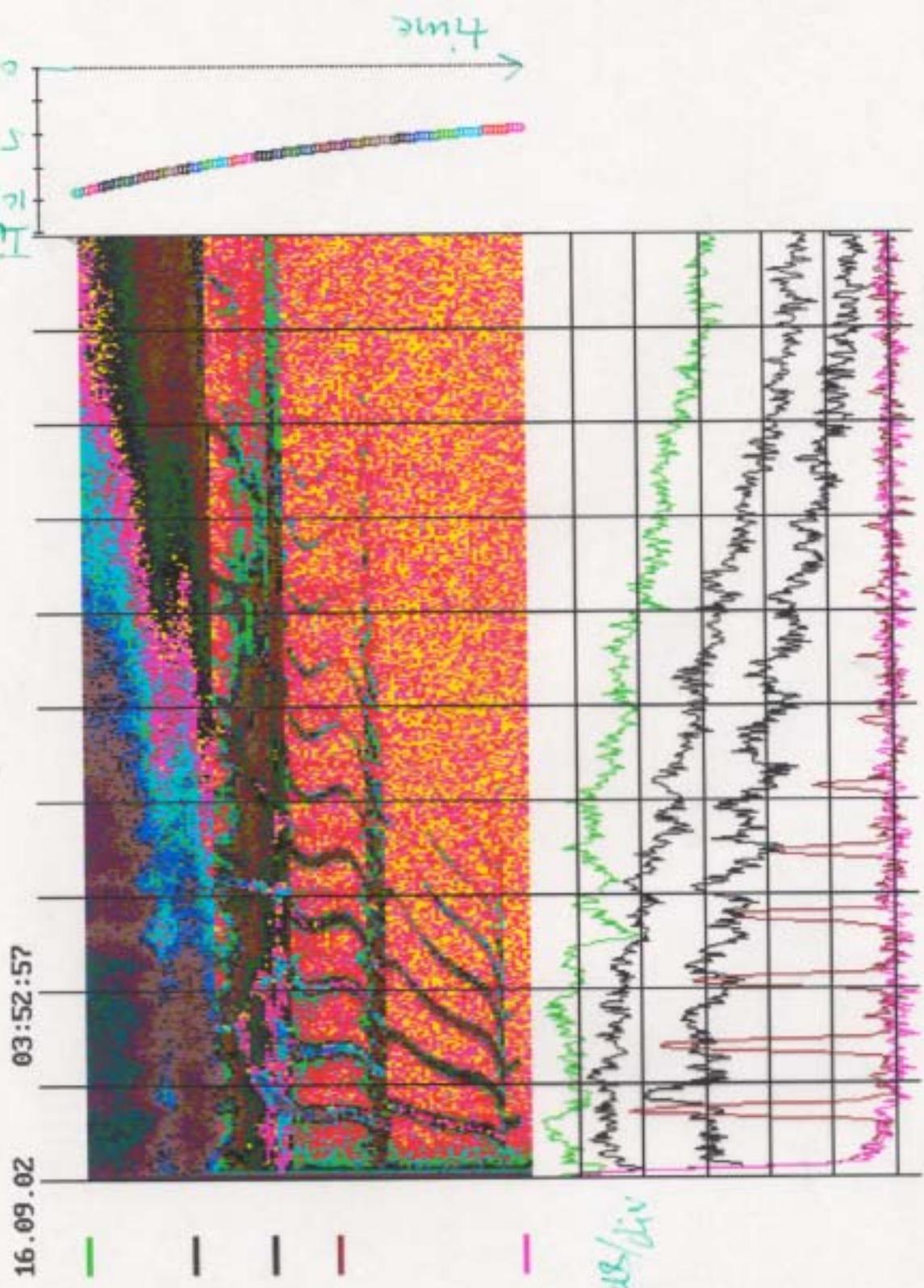
1000/div

20 kHz

10 kHz

c

Bursting CSR - two Runches



16.09.02 03:52:57

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-
-
-

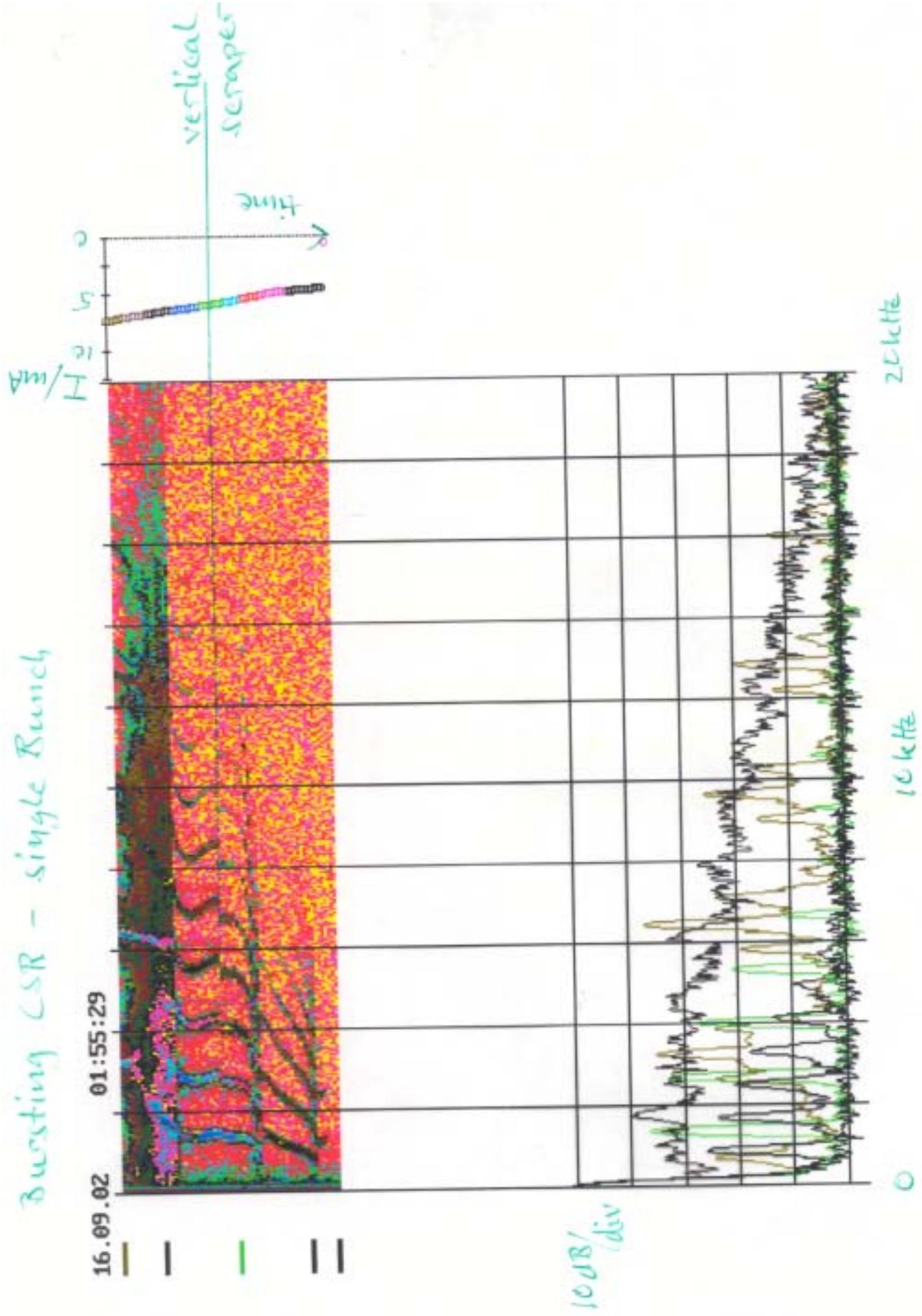
10 dB/div

20 kHz

10 kHz

0

Bursting CSR - single Bunch



CONCLUSION

Streak camera resolution limited by phase noise to:
Loss factors indicate much shorter bunches.

Bunch asymmetry explains extended wavelength range for stable CSR.

Contributions of k to the static non-linearity of the longitudinal single particle motion have been determined. Effect on CSR is small.

———— Above threshold ————

Energy spread increases

Quadrupole and higher azimuthal modes appear in the longitudinal beam spectra.

Emission of CSR in bursts. Above the threshold many transitions occur between regular and irregular repetition rate.

Modelling this behaviour requires to take into account CSR as well as chamber impedance